HARP-CDP hadroproduction data: Comparison with FLUKA and GEANT4 simulations

Abstract

We report on the comparison of production characteristics of secondary protons and charged pions in the interactions of protons and charged pions with momentum between 3 GeV/*c* and 15 GeV/*c* with beryllium, copper, and tantalum nuclei, with simulations by the FLUKA and Geant4 Monte Carlo tool kits. Overall production cross-sections are reasonably well reproduced, within factors of two. In more detail, there are areas with poor agreement that are unsatisfactory and call for modelling improvements. Overall, the current FLUKA simulation fares better than the current Geant4 simulation.

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1 INTRODUCTION

The HARP experiment arose from the realization that the inclusive differential cross-sections of hadron production in the interactions of few GeV/*c* protons with nuclei were known only within a factor of two to three, while more precise cross-sections are in demand for several reasons. Among them are the understanding of the underlying physics and the modelling of Monte Carlo generators of hadron–nucleus collisions.

The HARP experiment was designed to carry out a programme of systematic and precise (i.e., at the few per cent level) measurements of hadron production by protons and pions with momenta from 1.5 to 15 GeV/c, on a variety of target nuclei ranging from hydrogen to lead.

The HARP detector combined a forward spectrometer with a large-angle spectrometer. The latter comprised a cylindrical Time Projection Chamber (TPC) around the target and an array of Resistive Plate Chambers (RPCs) that surrounded the TPC. The purpose of the TPC was track reconstruction and particle identification by dE/dx. The purpose of the RPCs was to complement the particle identification by time of flight.

The HARP experiment took data at the CERN Proton Synchrotron in 2001 and 2002.

Several papers reported on the measurement of inclusive cross-sections of large-angle production (polar angle θ in the range $20^{\circ} < \theta < 125^{\circ}$) of secondary protons and charged pions, in the interactions with 5% λ_{int} beryllium, copper, tantalum and lead targets of protons and pions with beam momenta of $\pm 3.0, \pm 5.0, \pm 8.0$ (+8.9 for beryllium), ± 12.0 , and ± 15.0 GeV/c [1–5]. Besides, one paper [6] reported on disagreements between data and simulations by the Geant4 Monte Carlo tool kit [7], which led to significant improvements in its simulation code [8].

In this paper, we report in more detail on comparisons of hadroproduction data with simulations by the FLUKA [9] and Geant4 Monte Carlo tool kits.

Our work involves only the HARP large-angle spectrometer. The data analysis is based on our calibrations of the HARP TPC and RPCs published in Refs. [10, 11].

2 THE BEAMS AND THE HARP SPECTROMETER

The protons and pions were delivered by the T9 beam line in the East Hall of CERN's Proton Synchrotron. This beam line supports beam momenta between 1.5 and 15 GeV/c, with a momentum bite $\Delta p/p \sim 1\%$.

The beam instrumentation, the definition of the beam particle trajectory, the cuts to select 'good' beam particles, and the muon and electron contaminations of the particle beams, are the same as described, e.g., in Ref. [4]. The targets were discs made of high-purity material, 5% λ_{int} thick.

The momentum resolution $\sigma(1/p_T)$ of the HARP–TPC is typically 0.2 (GeV/c)⁻¹ and worsens towards small relative particle velocity β and small polar angle θ . The absolute momentum scale is determined to be correct to better than 2%, both for positively and negatively charged particles.

The polar angle θ is measured in the TPC with a resolution of ~13 mrad, for a proton with $p_{\rm T} = 500 \text{ MeV}/c$ in the TPC gas and a polar angle of $\theta = 60^{\circ}$. The polar-angle scale is correct to better than 2 mrad.

The TPC measures dE/dx with a resolution of 16% for a track length of 300 mm.

The intrinsic efficiency of the RPCs that surround the TPC is better than 98%.

The intrinsic time resolution of the RPCs is 127 ps and the system time-of-flight resolution (that includes the jitter of the arrival time of the beam particle at the target) is 175 ps.

To separate measured particles into species, we assign on the basis of dE/dx and β to each particle a probability of being a proton, a pion (muon), or an electron, respectively. The

probabilities add up to unity, so that the number of particles is conserved. These probabilities are used for weighting when entering tracks into plots or tables.

3 MODUS OPERANDI OF THE COMPARISON BETWEEN DATA AND SIMULATIONS

All data shown in this paper have been published [1–5] in the form of double-differential inclusive cross-sections $d^2\sigma/dpd\Omega$ [mbarn/(sr·GeV/c)], in the transverse-momentum range $0.10 < p_T < 1.25$ GeV/c and the polar-angle range $20^\circ < \theta < 125^\circ$. For ease of use, the cross-sections are also available in computer-readable form as ASCII files [12–16].

For the comparison with simulations, cross-sections are integrated over two regions: the 'intermediate-angle' region $(20^{\circ} < \theta < 50^{\circ})$ and the 'large-angle' region $(50^{\circ} < \theta < 125^{\circ})$. For π^+ and π^- secondaries, the $p_{\rm T}$ ranges are $0.10 < p_{\rm T} < 0.72$ GeV/c in the intermediateangle region, and $0.16 < p_{\rm T} < 1.25$ GeV/c in the large-angle region. For proton secondaries, the $p_{\rm T}$ range is $0.30 < p_{\rm T} < 0.72$ GeV/c, for the large energy loss by ionization in the target and in materials before entering the active TPC volume. Cross-sections of protons are given only in the intermediate-angle region because the minimum $p_{\rm T}$ of protons in the large-angle region is even larger, about twice the one in the intermediate-angle region.

The measured differential cross-sections in the said regions were obtained by integrating over respective bins of polar angle and transverse momentum, taking into account the correlation of systematic errors. The uncertainties of the shown inclusive cross-sections are at the level of 3%. The contribution of statistical errors is negligible, with the exception of the cross-sections for 15 GeV/c π^+ beams. The largest contributions to the systematic error arise from the overall normalization (2%) and from the uncertainty of the momentum scale of secondaries (2%).

The Monte Carlo tool kits FLUKA and Geant4¹⁾ are run with protons and charged-pion beams with the same beam momenta that interact with beryllium, copper and tantalum target nuclei. Only final-state hadrons that stem from the target nuclei are taken into account. For the same regions selected for the presentation of data, integrated inclusive cross-sections are extracted.

The comparison is made separately (i) for the intermediate-angle and the large-angle regions, (ii) for secondary protons, π^+ 's and π^- 's, and (iii) for incoming beam protons, π^+ 's and π^- 's in the momentum range 3–15 GeV/c.

For Geant4, the QGSP_BERT 'physics list' was selected for being the preferred choice of the LHC Collaborations ATLAS and CMS [17]. In order to assess the differences with other popular physics lists, in Fig. 1 QGSP_BERT simulations of inclusive π^+ production by protons on beryllium, copper and tantalum nuclei are compared with the simulations employing the Geant4 physics lists QGSP_BIC, QGSP_EMV and FTFP_BERT. There are remarkable differences between the simulations, similar in size as the differences between the data and, e.g., the QGSP_BERT simulation. Most of these differences show up for beam proton momenta below 10 GeV/c. Between 10 and 15 GeV/c beam proton momentum, the chosen four physics lists give comparable results.

Our aim is not to test all possible Geant4 physics lists against our data. Rather, we wish to point to areas, using the QGSP_BERT physics list as an example, where data and simulation seriously disagree.

We trust that our data are useful to the simulation developers and provide guidance for modelling improvements.

¹⁾We used the program versions FLUKA 2008.3c and Geant 4.9.3.



Fig. 1: Comparison of measured inclusive π^+ production cross-sections by protons on Be (squares), Cu (triangles) and Ta (circles) with Geant4 simulations employing the QGSP_BERT (upper left panel), QGSP_BIC (upper right panel), QGSP_EMV (lower left panel) and FTFP_BERT (lower right panel), in the intermediate-angle region, as a function of beam momentum; the data are shown with black symbols, the simulations with open symbols; here and in further similar figures, the simulated points are connected by lines to guide the eye.

The comparison of data with FLUKA and Geant4 QGSP_BERT simulations is concentrated in Section 4. A critical appraisal is found in Section 6.

4 COMPARISON OF DATA WITH FLUKA AND GEANT4 SIMULATIONS

Figures 2–4 show comparisons with FLUKA and Geant4 simulations of measured inclusive proton, π^+ and π^- production cross-sections by protons on Be, Cu and Ta nuclei. For final-state protons, only comparisons in the intermediate-angle region are given (Fig. 2). For final-state π^+ 's and π^- 's, comparisons are given in the intermediate-angle (Fig. 3) and large-angle regions (Fig. 4).

While Figs. 2–4 show comparisons for proton beam particles, Figures 5–7 and Figs. 8–10 show the same comparisons for π^+ and π^- beam particles, respectively.



Fig. 2: Comparison of measured (black symbols) inclusive proton production cross-sections by protons on Be, Cu and Ta nuclei, in the intermediate-angle region, as a function of beam momentum, with FLUKA (upper panel) and Geant4 (lower panel) simulations (open symbols).



Fig. 3: Comparison of measured inclusive π^+ (left panels) and π^- (right panels) production crosssections by protons on Be, Cu and Ta nuclei (black symbols), in the intermediate-angle region, with FLUKA and Geant4 simulations (open symbols).



Fig. 4: Same as Fig. 3 but in the large-angle region.



Fig. 5: Comparison of measured (black symbols) inclusive proton production cross-sections by π^+ 's on Be, Cu and Ta nuclei, in the intermediate-angle region, as a function of beam momentum, with FLUKA (upper panel) and Geant4 (lower panel) simulations (open symbols).



Fig. 6: Comparison of measured inclusive π^+ (left panels) and π^- (right panels)production crosssections by π^+ 's on Be, Cu and Ta nuclei (black symbols), in the intermediate-angle region, with FLUKA and Geant4 simulations (open symbols).



Fig. 7: Same as Fig. 6 but in the large-angle region.



Fig. 8: Comparison of measured (black symbols) inclusive proton production cross-sections by π^- 's on Be, Cu and Ta nuclei, in the intermediate-angle region, as a function of beam momentum, with FLUKA (upper panel) and Geant4 (lower panel) simulations (open symbols).



Fig. 9: Comparison of measured inclusive π^+ (left panels) and π^- (right panels)production crosssections by π^- 's on Be, Cu and Ta nuclei (black symbols), in the intermediate-angle region, with FLUKA and Geant4 simulations (open symbols).



Fig. 10: Same as Fig. 9 but in the large-angle region.

5 EFFECTS FROM RE-INTERACTIONS OF HADRONS IN NUCLEAR MATTER

Before our appraisal of the comparison of data with simulations in Section 6, it is useful to point to generic features of final-state hadrons from heavy nuclei.

Fig. 11 shows inclusive cross-sections of proton-production with polar angle $20^{\circ} < \theta < 50^{\circ}$ by incoming π^- beam particles with momenta in the range 3–15 GeV/*c*, for beryllium and and tantalum nuclei. The tantalum cross-sections are shown as measured, while the beryllium cross-sections are scaled with $(A_{\rm Ta}/A_{\rm Be})^{0.7}$. The rationale is that the latter cross-sections are the ones on a hypothetical tantalum nucleus that has the same small re-interaction probability of secondaries as a beryllium nucleus. Fig. 11 tells that the measured inclusive cross-section of tantalum is still four times larger than the scaled beryllium cross-section. Therefore, the bulk of final-state protons stem from re-interactions of secondaries in the nuclear matter of the tantalum nucleus—which after all has a diameter that is equivalent to several nuclear interaction lengths.



Fig. 11: Inclusive cross-sections of proton-production with polar angle $20^{\circ} < \theta < 50^{\circ}$ by incoming π^{-} beam particles with momenta in the range 3–15 GeV/*c*, for beryllium and and tantalum nuclei; The tantalum cross-sections are those measured, while the beryllium cross-sections are scaled with $(A_{\rm Ta}/A_{\rm Be})^{0.7}$.

In Fig. 11, it also appears that the inclusive cross-section of proton production on tantalum nuclei at 3 GeV/*c* beam momentum does not follow the trend that is suggested by the data points at higher beam momentum. Indeed, a similar but smaller reduction of cross-section can be seen in the proton production by incoming beam π^+ 's (see Fig. 5) and an even smaller reduction also for incoming beam protons (see Fig. 2). No such reduction is seen for the production of secondary π^+ 's and π^- 's. We conjecture that this reduction reflects the absorption of low-momentum final-state protons in the nuclear matter of the target nucleus.

With a view to corroborating this conjecture, Fig. 12 shows the measured total-momentum distributions of final-state protons, in the polar-angle range $20^{\circ} < \theta < 50^{\circ}$, for incoming π^{-} beam particles. Since the emphasis is on shape comparison, all distributions are normalized to unit area. The upper panel shows the comparison between beam momenta of 3, 8 and 15 GeV/*c* beam momentum for beryllium target nuclei, while the lower panel shows the same for tantalum target nuclei. The difference in the shape of the total-momentum distributions of final-state



Fig. 12: Comparison of total-momentum distributions, normalized to unit area, of final-state protons with polar angle $20^{\circ} < \theta < 50^{\circ}$, for π^{-} beam particle momenta of 3 GeV/*c* (black points), 8 GeV/*c* (dotted histogram) and 15 GeV/*c* (full histogram), for beryllium (upper panel) and for tantalum (lower panel) target nuclei.

protons between Be and Ta, and for Ta between different π^- beam particle momenta, is apparent. The heavier the nucleus and the lower the beam momentum, the more the spectrum is shifted toward smaller momentum values. Because of the strong rise of the proton–nucleon cross-section for proton momenta below 0.8 GeV/*c*, many final-state protons are absorbed inside the target nucleus which causes a cross-section reduction which affects among our beam momenta strongest the final-state protons at 3 GeV/*c* beam momentum. The reduction is largest for π^- beam particles because they are least efficient to produce high-momentum final-state protons that escape the nucleus. The reduction is not seen for final-state pions because the pion–nucleon cross-section permits at low momenta for pions an easier escape from the nucleus than for protons.

6 APPRAISAL OF THE DIFFERENCES BETWEEN DATA AND SIMULATIONS

6.1 FLUKA versus data

In our parameter range of the comparison of data with FLUKA simulations, the latter agree with data within some 30%.

The most unsatisfactory feature in the FLUKA simulation is the discontinuity around 5 GeV/*c* beam momentum which is rather persistent in all combinations of beam particle and final-state particle.

6.2 Geant4 versus data

In our parameter range of the comparison of data with Geant4 QGSP_BERT simulations, the latter agree with data within a factor of about two.

An unacceptable feature of the Geant4 QGSP_BERT simulation is the strong discontinuity around 10 GeV/c beam momentum which is also rather persistent in all combinations of beam particle and final-state particle.

With a view to exploring further this discontinuity around 10 GeV/c in Geant4, Fig. 13 shows the simulation of the inclusive generation of secondary neutrons and protons with a



Fig. 13: Geant4 simulation (QGSP_BERT physics list) of the generation of secondary protons and neutrons with $p_{\rm T} > 0.1$ GeV/c in the interaction of protons with Ta nuclei.

rather loose cut of $p_{\rm T} > 0.1$ GeV/c, in the interaction of beam protons with momentum in the range 3–30 GeV/c with Ta nuclei. Most prominent is an abrupt and unphysical change of modelling around 10 GeV/c. Further, a comparison with the large proton multiplicity shown in Fig. 13—up to some 20 protons per beam particle interaction—with the measured much smaller proton multiplicity (see, e.g., Ref. [5]) suggests that final-state protons are dominated by protons from the fragmentation of the target nucleus and not by final-state protons from the interactions of the primary beam particle or of secondary hadrons in the nuclear matter of the target nucleus. Since for protons from the fragmentation of the target nucleus it is expected that the momentum spectrum of secondary protons is steeply falling with increasing momentum, the cut $p_{\rm T} > 0.3$ GeV/c will strongly diminish the protons from the fragmentation of the target nucleus.

Fig. 14 shows the polar-angle distribution of protons generated by Geant4 (QGSP_BERT physics list), in the interaction of protons with Ta nuclei. The open histograms refer to proton beam momenta below the discontinuity around 10 GeV/c beam momentum, the shaded histograms to proton beam momenta above. The upper panel has no cut applied while the lower panel shows protons with $p_{\rm T} > 0.3$ GeV/c. The discrepancy between the distributions below and above the discontinuity around 10 GeV/c is unphysical.

Figure 15 shows the polar-angle distributions of π^+ 's generated by protons interacting with Ta nuclei. No cuts are applied. The upper panel shows the FLUKA simulation, the lower panel the Geant4 (QGSP_BERT physics list) simulation. Here the open histograms refer to proton beam momenta just above the discontinuities discussed before (around 5 GeV/c for FLUKA, around 10 GeV/c for Geant4), the shaded histograms to proton beam momenta just below. The differences of the polar-angle distributions below and above the discontinuities are unphysical.

Based on the same raw data from which we extracted our inclusive cross-sections, the HARP Collaboration published inclusive cross-sections of the production of π^+ and π^- (but not of protons) by incoming proton, π^+ and π^- beam particles, in the polar-angle region 0.35 rad $\leq \theta \leq 2.15$ rad [18,19], and compared their results with predictions of the Geant4 and MARS [20] Monte Carlo tool kits. Substantial discrepancies were observed between their and our results, documented in Refs. [1–5], discussed in Refs [21–25], and summarized in the Appendix of Ref. [1].



Fig. 14: The polar-angle distribution of protons generated by Geant 4 (QGSP_BERT physics list) in the interaction of protons with Ta nuclei; the open histograms refer to proton beam momenta just below the discontinuity around 10 GeV/c beam momentum, the shaded (yellow) histograms to proton beam momenta just above; the upper panel has no cut applied while the lower panel shows protons with $p_{\rm T} > 0.3 \text{ GeV/}c$.



Fig. 15: The polar-angle distributions of π^+ 's generated by protons interacting with Ta nuclei, with no cuts applied; the upper panel shows the FLUKA simulation, the lower panel the Geant4 (QGSP_BERT physics list) simulation; the open histograms refer to proton beam momenta just above the discontinuities (around 5 GeV/*c* for FLUKA, around 10 GeV/*c* for Geant4), the shaded (yellow) histograms to proton beam momenta just below.

7 SYNOPSIS

A comprehensive comparison of measured inclusive cross-sections of proton, π^+ and π^- production by beams of protons, π^+ 's and π^- 's with momentum in the range 3–15 GeV/*c*, interacting with beryllium, copper and tantalum target nuclei, with simulations by the FLUKA and Geant4 Monte Carlo tool kits is presented. Overall production cross-sections are reasonably well reproduced, within factors of two. In more detail, there are areas with poor agreement that are unsatisfactory and call for modelling improvements. Overall, the current FLUKA simulation fares better than the current Geant4 simulation.

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