

TOTAL DESTRUCTION OF LAGRANGIAN TORI

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ABSTRACT. For an integrable Tonelli Hamiltonian with d ($d \geq 2$) degrees of freedom, we show that all of the Lagrangian tori can be destroyed by analytic perturbations which are arbitrarily small in the $C^{d-\delta}$ topology.

Key words. Lagrangian torus, Tonelli Hamiltonian

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1. INTRODUCTION AND MAIN RESULT

For exact area-preserving twist maps on annulus, it was proved by Herman in [H2] that invariant circles with given rotation numbers can be destroyed by $C^{3-\delta}$ arbitrarily small C^∞ perturbations, where δ is a small positive constant. Following the ideas and techniques developed by J.N.Mather in the series of papers [M1], [M2], [M3] and [M4], a variational proof of Herman's result was provided in [Wa]. A Hamiltonian is called a Tonelli Hamiltonian if it satisfies positive definiteness, superlinear growth with respect to momentum and completeness of flow. For the case with $d \geq 2$ degrees of freedom, it was proved in [CW] that for every given rotation vector, invariant Lagrangian torus with that rotation vector of an integrable Tonelli Hamiltonian system can be destroyed by an arbitrarily small C^∞ perturbation in the $C^{2d-\delta}$ topology. In contrast with it, it was shown that KAM torus with Diophantine frequency persists under $C^{2d+\delta}$ -small perturbations ([Po]). Hence, the above result is almost optimal.

On the other hand, it was proved that all invariant circles can be destroyed by C^1 arbitrarily small C^∞ perturbations of the integrable area-preserving twist maps [Ta]. C^1 topology was improved to be $C^{2-\delta}$ topology by Herman in [H1]. Moreover, he extended the result to systems with multi-degrees of freedom and found that all of the Lagrangian tori of an integrable symplectic twist map can be destroyed by $C^{d+2-\delta}$ arbitrarily small C^∞ perturbations of the generating function [H3]. By [H2], each orbit on an invariant Lagrangian graphs is an action minimizing curve. Based on the minimality of the orbits on the Lagrangian graph, [MMS] provided a criterion of non-existence of all Lagrangian tori and applied it to a particular example. Based on the correspondence between symplectic twist maps and Tonelli Hamiltonian systems ([Go, Mo]), it shows that all of the Lagrangian tori of an integrable Tonelli Hamiltonian system with $d \geq 2$ degrees of freedom can be destroyed by C^∞ perturbations which are arbitrarily small in the $C^{d+1-\delta}$ topology.

Comparing the results on both sides, it is natural to ask the following question:

- if all of the Lagrangian tori can be destroyed by an arbitrarily small real-analytic perturbation in the C^r topology, what is the maximum of r ?

In this paper, we prove the following theorem:

Theorem 1.1 *For an integrable Tonelli Hamiltonian with d ($d \geq 2$) degrees of freedom, all of the Lagrangian tori can be destroyed by analytic perturbations which are arbitrarily small in the $C^{d-\delta}$ topology for a small given constant $\delta > 0$.*

Unfortunately, we still don't know whether the other results are optimal. Some further developments of KAM theory are needed to verify the optimality. To prove Theorem 1.1, we recall some notions on Lagrangian torus.

In $T^*\mathbb{T}^d$, a submanifold \mathcal{T}^d is called Lagrangian torus if it is diffeomorphic to the torus \mathbb{T}^d and the symplectic form (non-degenerate closed 2-form) vanishes on it. An example of Lagrangian torus is the KAM torus.

Definition 1.2 $\bar{\mathcal{T}}^d$ is called a d dimensional KAM torus if

- $\bar{\mathcal{T}}^d$ is a Lipschitz graph over \mathbb{T}^d ;
- $\bar{\mathcal{T}}^d$ is invariant under the Hamiltonian flow Φ_t^H generated by the Hamiltonian function H ;
- there exists a diffeomorphism $\phi : \mathbb{T}^d \rightarrow \bar{\mathcal{T}}^d$ such that $\phi^{-1} \circ \Phi_H^t \circ \phi = R_\omega^t$ for any $t \in \mathbb{R}$, where $R_\omega^t : x \rightarrow x + \omega t$ and ω is called the rotation vector of $\bar{\mathcal{T}}^d$.

Generally, the rotation vector of the Lagrangian torus is not well defined if the Lagrangian torus contains several invariant sets with different rotation vectors. In this paper, we are concerned with the Lagrangian torus as follow:

Definition 1.3 \mathcal{T}^d is called a d dimensional Lagrangian torus if

- \mathcal{T}^d is a Lagrangian graph;
- \mathcal{T}^d is invariant for the Hamiltonian flow Φ_H^t generated by H .

It is still open whether all the invariant Lagrangian tori of Tonelli Hamiltonians are graphs or not. Some results have been obtained by adding topological or variational conditions (see [Ar], [BP] and [CR] for instance). [Ar] proved that for Tonelli Hamiltonians, the invariant Lagrangian submanifolds that are isotopic to zero section are graphs. It is easy to see that "isotopic to zero section" is necessary if we consider the elliptic island for the Hamiltonian time one map of rigid pendulum. However, the orbits on the elliptic island is not minimal in the sense of variation. Under the condition that every orbit is an action minimizing curve, [CR] proved that for Tonelli Hamiltonians with 2 degrees of freedom, the invariant Lagrangian tori are graphs. By [H2], each orbit on an invariant Lagrangian graph is an action minimizing curve. Hence, it is shown that the graph property of a Lagrangian torus is equivalent to the minimality of the orbits on the torus for the case with two degrees of freedom. Whereas, it is still open to verify the equivalence for the case with multi-degrees of freedom.

In the following sections, we consider the destruction of all of the Lagrangian tori of symplectic twist maps.

Definition 1.4 A map $f : (x, y) \rightarrow (x', y')$ of $\mathbb{T}^d \times \mathbb{R}^d$ is called a symplectic twist map if

- f is a diffeomorphism isotopic to the identity;
- f preserves the symplectic form;

- $\frac{\partial x'}{\partial y}$ is uniformly positive definite and bounded.

In particular, the last one is called twist condition. An necessary condition for existence of any invariant Lagrangian graph is that there exists a generating function $h : \mathbb{T}^d \times \mathbb{R}^d \rightarrow \mathbb{R}$ such that

$$y' = \partial_2 h(x, x') \quad y = -\partial_1 h(x, x')$$

where ∂_1, ∂_2 denote derivatives with respect to the first and second arguments respectively. We will restrict attention to maps with periodic generating functions, also known as exact symplectic maps.

Based on the correspondence between symplectic twist maps and Hamiltonian systems, it can be achieved to destroy all of the Lagrangian tori of Hamiltonian systems.

2. A TOY MODEL

To show the basic ideas, we are beginning with a toy model whose generating function is as follow:

$$(2.1) \quad h_n(x, x') = h_0(x, x') - \frac{5}{4n^2} \sin(nx') - \frac{1}{16n^2} \cos(2nx'),$$

where $h_0(x, x') = \frac{1}{2}(x - x')^2$. Let $f_n(x, y) = (x', y')$ be the exact area-preserving twist map generated by (2.1), then

$$\begin{cases} y = -\partial_1 h_n(x, x') = x' - x, \\ y' = \partial_2 h_n(x, x') = x' - x - \frac{5}{4n} \cos(nx') + \frac{1}{8n} \sin(2nx'). \end{cases}$$

We set $\phi_n(x) = -\frac{5}{4n} \cos(nx) + \frac{1}{8n} \sin(2nx)$, then

$$(2.2) \quad f_n(x, y) = (x + y, y + \phi_n(x + y)).$$

In [H1], Herman found a criterion of total destruction of invariant circles. By Birkhoff graph theorem (see [H2]), if f_n admits an invariant circle, then the invariant circle is a Lipschitz graph. We denote the graph by ψ_n , then it follows from [H3] that

$$\psi_n \circ g_n - \psi_n = \phi_n \circ g_n,$$

where $g_n = \text{Id} + \psi_n$. This is equivalent to

$$(2.3) \quad \frac{1}{2}(g_n + g_n^{-1}) = \text{Id} + \frac{1}{2}\phi_n.$$

Let \mathfrak{D}_n be the set of differentiable points of g_n , then \mathfrak{D}_n has full Lebesgue measure on \mathbb{R} since g_n is a Lipschitz function. For $x \in \mathfrak{D}_n$, we differentiate (2.3),

$$\frac{1}{2}(Dg_n(x) + (Dg_n)^{-1}(g_n^{-1}(x))) = 1 + \frac{1}{2}D\phi_n(x).$$

Let $G_n = \|Dg_n\|_{L^\infty}$. It is easy to see that for $\varepsilon > 0$, there exists $\tilde{x} \in \mathfrak{D}_n$ such that $Dg_n(\tilde{x}) \geq G_n - \varepsilon$. Let $M_n = \max D\phi_n$, we have

$$\frac{1}{2} \left(G_n + \frac{1}{G_n} - \varepsilon \right) \leq 1 + \frac{1}{2}M_n.$$

Since $\varepsilon > 0$ is arbitrarily small, then

$$\frac{1}{2} \left(G_n + \frac{1}{G_n} \right) \leq 1 + \frac{1}{2} M_n.$$

Hence,

$$(2.4) \quad G_n \leq 1 + \frac{1}{2} M_n + \left(M_n + \frac{1}{4} M_n^2 \right)^{\frac{1}{2}}.$$

Obviously, for $x \in \mathfrak{D}_n$, we have

$$\frac{1}{G_n} \leq 1 + \frac{1}{2} D\phi_n(x).$$

Let $m_n = \min D\phi_n$, then we have

$$\frac{1}{G_n} \leq 1 + \frac{1}{2} m_n,$$

which together with (2.4) implies that

$$\frac{1}{1 + \frac{1}{2} m_n} \leq 1 + \frac{1}{2} M_n + \left(M_n + \frac{1}{4} M_n^2 \right)^{\frac{1}{2}}.$$

Therefore, it is sufficient for total destruction of invariant circles to construct $\phi_n(x)$ such that

$$(2.5) \quad \frac{1}{1 + \frac{1}{2} \min D\phi_n} > 1 + \frac{1}{2} \max D\phi_n + \left(\max D\phi_n + \frac{1}{4} (\max D\phi_n)^2 \right)^{\frac{1}{2}}.$$

In our construction,

$$(2.6) \quad D\phi_n(x) = \frac{5}{4} \sin(nx) + \frac{1}{4} \cos(2nx).$$

A simple calculation implies

$$\begin{cases} \min D\phi_n(x) = -\frac{3}{2}, & \text{attained at } x = \frac{3}{2n} + \frac{2\pi k}{n}, \\ \max D\phi_n(x) = 1, & \text{attained at } x = \frac{\pi}{2n} + \frac{2\pi k}{n}, \end{cases}$$

where $k \in \mathbb{Z}$. Hence, (2.5) holds. Moreover, the exact area-preserving twist map generated by (2.1) admits no invariant circles.

By interpolation inequality ([H1]), for a small positive constant δ , we have

$$\|\phi_n\|_{C^{1-\delta}} \leq 2 \|\phi_n\|_{C^0}^\delta \|D\phi_n\|_{C^0}^{1-\delta}.$$

From the construction of ϕ_n , it follows that $\|\phi_n\|_{C^0} \rightarrow 0$, as $n \rightarrow \infty$ and $\|D\phi_n\|_{C^0}$ is bounded. Hence,

$$\|\phi_n\|_{C^{1-\delta}} \rightarrow 0 \quad \text{as } n \rightarrow \infty,$$

which implies that

$$\|h_n - h_0\|_{C^{2-\delta}} \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

3. C^∞ DESTRUCTION OF ALL OF THE LAGRANGIAN TORI

In [H3], Herman extended the criterion (2.5) to multi-degrees of freedom. More precisely, for exact symplectic twist map on $\mathbb{T}^{*r}\mathbb{T}^d$, whose generating function is

$$(3.1) \quad h(x, x') = \frac{1}{2}(x - x')^2 + \Psi(x'),$$

where $\Psi \in C^r(\mathbb{T}^d, \mathbb{R})$, $r \geq 2$. Correspondingly, the exact symplectic twist map has the following form

$$(3.2) \quad f(x, y) = (x + y, y + d\Psi(x + y)),$$

where

$$d\Psi = \left(\frac{\partial \Psi}{\partial x_1}, \dots, \frac{\partial \Psi}{\partial x_d} \right).$$

Let $E(x)$ be the derivative matrix of $d\Psi$ and $T(x) = \frac{1}{d}\text{tr}E(x)$, where $\text{tr}E(x)$ denotes the trace of $E(x)$. From a similar argument as the deduction of (2.5), it follows that it is sufficient for total destruction of the Lagrangian tori to construct $T(x)$ such that

$$(3.3) \quad \frac{1}{1 + \frac{1}{2}\min T(x)} > 1 + \frac{1}{2}\max T(x) + \left(\max T(x) + \frac{1}{4}(\max T(x))^2 \right)^{\frac{1}{2}}.$$

Moreover, for $T(x) \rightarrow 0$, (3.3) implies

$$(3.4) \quad -\frac{1}{2}\min T(x) > \sqrt{\max T(x)} + O(\max T(x)).$$

Remark 3.1 *The integrable part $\frac{1}{2}(x - x')^2$ can be easily generalized to the form $\frac{1}{2}(x - x')^t A(x - x')$, where $(\cdot)^t$ denotes the transpose of (\cdot) and A denotes a symmetric positive definite matrix. By a similar calculation, (3.3) still holds true if $T(x) = \frac{1}{d}\text{tr}E(x)$ is replaced by $T(x) = \frac{1}{d}\text{tr}(A^{-1}E(x))$. The choice of the form $\frac{1}{2}(x - x')^2$ could set us free from a tedious calculation to see the crucial mechanism of the problem.*

Herman constructed a sequence $\{\Psi_n\}_{n \in \mathbb{N}}$ that satisfies (3.4). It is easy to see $T_n(x) = \frac{1}{d}\Delta\Psi_n$ where Δ denotes the Laplacian. Since $T_n(x)$ is 2π -periodic, it is enough to construct it on $[-\pi, \pi]^d$. More precisely,

$$T_n(x) = \begin{cases} T_n^+(x), & x \in [0, \pi]^d, \\ -T_n^-(x), & x \in [-\pi, 0]^d, \\ 0, & \text{others.} \end{cases}$$

where $T_n(x)$ is C^∞ function, $T_n^+(x)$ and $T_n^-(x)$ have the following forms respectively.

$T_n^+(x)$ satisfies:

$$\begin{cases} \text{supp } T_n^+(x) \subset [0, \pi]^d, \\ \max T_n^+(x) = \frac{1}{9n}, \end{cases}$$

$T_n^-(x)$ satisfies:

$$\begin{cases} \text{supp } T_n^-(x) = B_{R_n}(x_0), \\ \max T_n^-(x) = \frac{1}{\sqrt{n}}, \\ R_n \sim \left(\frac{1}{\sqrt{n}} \right)^{\frac{1}{d}}, \\ x_0 = \left(-\frac{\pi}{2}, \dots, -\frac{\pi}{2} \right), \end{cases}$$

where $f \sim g$ means that $\frac{1}{C}g < f < Cg$ holds for a constant $C > 1$. Hence, we obtain a sequence of $\{T_n(x)\}_{n \in \mathbb{N}}$ with bounded C^d norms and satisfying

$$\int_{\mathbb{T}^d} T_n(x) dx = 0.$$

From interpolation inequality, it follows that $T_n(x) \rightarrow 0$ as $n \rightarrow \infty$ in the $C^{d-\delta}$ topology for any $\delta > 0$.

Let Ψ_n be the unique function in $C^\infty(\mathbb{T}^d, \mathbb{R})$ such that

$$\int_{\mathbb{T}^d} \Psi_n(x) dx = 0 \quad \text{and} \quad \frac{1}{d} \Delta \Psi_n(x) = T_n(x).$$

By Schauder estimates one knows that for any $\delta > 0$, $\Psi_n(x) \rightarrow 0$ as $n \rightarrow \infty$ in the $C^{d+2-\delta}$ topology. From the construction of $T_n(x)$, it is easy to see that (3.4) is verified.

Above all, we have the following theorem

Theorem 3.2 *All of the Lagrangian tori of an integrable symplectic twist map with $d \geq 1$ degrees of freedom can be destroyed by C^∞ perturbations of the generating function and the perturbations are arbitrarily small in the $C^{d+2-\delta}$ topology for a small given constant $\delta > 0$.*

Based on the correspondence between symplectic twist maps and Hamiltonian systems, we have the following corollary.

Corollary 3.3 *All of the Lagrangian tori of an integrable Tonelli Hamiltonian system with $d \geq 2$ degrees of freedom can be destroyed by C^∞ perturbations which are arbitrarily small in the $C^{d+1-\delta}$ topology for a small given constant $\delta > 0$.*

4. AN APPROXIMATION LEMMA

In this section, we will prove a lemma on C^∞ functions approximated by trigonometric polynomials. First of all, we need some notations. Define

$$C_{2\pi}^\infty(\mathbb{R}^d, \mathbb{R}) := \left\{ f : \mathbb{R}^d \rightarrow \mathbb{R} \mid f \in C^\infty(\mathbb{R}^d, \mathbb{R}) \text{ and } 2\pi\text{-periodic in } x_1, \dots, x_d \right\}.$$

Let $f(x) \in C_{2\pi}^\infty(\mathbb{R}^d, \mathbb{R})$. The m -th Fejér-polynomial of f with respect to x_j is given by

$$(4.1) \quad F_m^{[j]}(f)(x) := \frac{1}{m\pi} \int_{-\pi/2}^{\pi/2} f(x + 2te_j) \left(\frac{\sin(mt)}{\sin t} \right)^2 dt,$$

where $x \in \mathbb{R}^d$, $m \in \mathbb{N}$, $j \in \{1, \dots, d\}$ and e_j is the j -th vector of the canonical basis of \mathbb{R}^d . $F_m^{[j]}(f)(x)$ is a trigonometric polynomial in x_j of degree at most $m - 1$. By [Zy],

$$\frac{1}{m\pi} \int_{-\pi/2}^{\pi/2} \left(\frac{\sin(mt)}{\sin t} \right)^2 dt = 1,$$

hence, from (4.1), we have

$$\|F_m^{[j]}(f)\|_{C^0} \leq \|f\|_{C^0}.$$

We denote

$$P_m^{[j]}(f) := 2F_{2m}^{[j]}(f) - F_m^{[j]}(f).$$

It is easy to see that $P_m^{[j]}(f)$ is a trigonometric polynomial in x_j of degree at most $2m - 1$. Moreover,

$$(4.2) \quad \|P_m^{[j]}(f)\|_{C^0} \leq 3\|f\|_{C^0},$$

$$(4.3) \quad P_m^{[j]}(af + bg) = aP_m^{[j]}(f) + bP_m^{[j]}(g),$$

where $a, b \in \mathbb{R}$ and $f, g \in C_{2\pi}^\infty(\mathbb{R}^d, \mathbb{R})$. For $k \in \{1, \dots, d\}$, $j_1, \dots, j_k \in \{1, \dots, d\}$ with $j_p \neq j_q$ for $p \neq q$. Let $m_1, \dots, m_k \in \mathbb{N}$ and $f \in C_{2\pi}^\infty(\mathbb{R}^d, \mathbb{R})$, we define

$$(4.4) \quad P_{m_1, \dots, m_k}^{[j_1, \dots, j_k]}(f) := P_{m_1}^{[j_1]} \left(P_{m_2}^{[j_2]} \left(\dots \left(P_{m_k}^{[j_k]}(f) \right) \dots \right) \right).$$

It is easy to see that for all $l \in \{1, \dots, k\}$, $P_{m_1, \dots, m_k}^{[j_1, \dots, j_k]}(f)$ are trigonometric polynomials in x_{j_l} of degree at most $2m_l - 1$, also known as generalized de la Vallée Poussin polynomial. We have the following lemma.

Lemma 4.1 *Let $f \in C_{2\pi}^\infty(\mathbb{R}^d, \mathbb{R})$, $r_1, \dots, r_d \in \mathbb{N}$, $m_1, \dots, m_d \in \mathbb{N}$, then we have*

$$(4.5) \quad \|f - P_{m_1, \dots, m_d}^{[1, \dots, d]}(f)\|_{C^0} \leq C_d \sum_{j=1}^d \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0},$$

where C_d is a constant only depending on d .

Lemma 4.1 is a direct corollary of Theorem 2.12 of [Al]. For the sake of completeness, we decide to provide the proof.

Proof We will prove Lemma 4.1 by induction. The case $d = 1$ is covered by Jackson's approximation theorem. More precisely, for $f \in C_{2\pi}^\infty(\mathbb{R}, \mathbb{R})$, $m, r \in \mathbb{N}$, we have

$$(4.6) \quad \|f - P_m^{[1]}(f)\|_{C^0} \leq C_1 \frac{1}{m^r} \left\| \frac{\partial^r f}{\partial x^r} \right\|_{C^0}.$$

Let the assertion be true for some $d \in \mathbb{N}$. We verify it for $d + 1$. Consider the functions $f(x_1, \cdot)$ with x_1 as a real parameter. Then by the assertion for d , we have

$$\|f(x_1, \cdot) - P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f)(x_1, \cdot)\|_{C^0} \leq C_d \sum_{j=2}^{d+1} \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0},$$

hence,

$$(4.7) \quad \|f - P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f)\|_{C^0} \leq C_d \sum_{j=2}^{d+1} \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0}.$$

Let $\hat{x}_j \in \mathbb{R}^d$ denote the vector $x \in \mathbb{R}^{d+1}$ without its j -th entry. For the functions $f(\cdot, \hat{x}_1)$, from (4.6), it follows that

$$\|f(\cdot, \hat{x}_1) - P_{m_1}^{[1]}(f)(\cdot, \hat{x}_1)\|_{C^0} \leq C_1 \frac{1}{m_1^{r_1}} \left\| \frac{\partial^{r_1} f}{\partial x_1^{r_1}} \right\|_{C^0},$$

hence,

$$(4.8) \quad \|f - P_{m_1}^{[1]}(f)\|_{C^0} \leq C_1 \frac{1}{m_1^{r_1}} \left\| \frac{\partial^{r_1} f}{\partial x_1^{r_1}} \right\|_{C^0}.$$

By (4.2), (4.3), (4.4) and (4.7), we have

$$\begin{aligned} \left\| P_{m_1}^{[1]}(f) - P_{m_1, \dots, m_{d+1}}^{[1, \dots, d+1]}(f) \right\|_{C^0} &= \left\| P_{m_1}^{[1]}(f) - P_{m_1}^{[1]} \left(P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f) \right) \right\|_{C^0}, \\ &= \left\| P_{m_1}^{[1]} \left(f - P_{m_1}^{[1]} P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f) \right) \right\|_{C^0}, \\ &\leq 3 \left\| f - P_{m_1}^{[1]} P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f) \right\|_{C^0}, \\ &\leq 3C_d \sum_{j=2}^{d+1} \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0}, \end{aligned}$$

which together with (4.8) implies that

$$\begin{aligned} \left\| f - P_{m_1, \dots, m_{d+1}}^{[1, \dots, d+1]}(f) \right\|_{C^0} &\leq \left\| f - P_{m_1}^{[1]}(f) \right\|_{C^0} + \left\| P_{m_1}^{[1]}(f) - P_{m_1, \dots, m_{d+1}}^{[1, \dots, d+1]}(f) \right\|_{C^0}, \\ &= \left\| f - P_{m_1}^{[1]}(f) \right\|_{C^0} + \left\| P_{m_1}^{[1]}(f) - P_{m_1}^{[1]} \left(P_{m_2, \dots, m_{d+1}}^{[2, \dots, d+1]}(f) \right) \right\|_{C^0}, \\ &\leq C_1 \frac{1}{m_1^{r_1}} \left\| \frac{\partial^{r_1} f}{\partial x_1^{r_1}} \right\|_{C^0} + 3C_d \sum_{j=2}^{d+1} \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0}, \\ &\leq C_{d+1} \sum_{j=1}^{d+1} \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0}. \end{aligned}$$

This finishes the proof of Lemma 4.1. \square

Obviously, there exist $m_{\bar{j}}, r_{\bar{j}}$ such that

$$\frac{1}{m_{\bar{j}}^{r_{\bar{j}}}} \left\| \frac{\partial^{r_{\bar{j}}} f}{\partial x_{\bar{j}}^{r_{\bar{j}}}} \right\|_{C^0} = \max_{1 \leq j \leq d} \left\{ \frac{1}{m_j^{r_j}} \left\| \frac{\partial^{r_j} f}{\partial x_j^{r_j}} \right\|_{C^0} \right\}.$$

Hence, we have

$$\begin{aligned} \|f - P_{m_1, \dots, m_d}^{[1, \dots, d]}(f)\|_{C^0} &\leq dC_d \frac{1}{m_{\bar{j}}^{r_{\bar{j}}}} \left\| \frac{\partial^{r_{\bar{j}}} f}{\partial x_{\bar{j}}^{r_{\bar{j}}}} \right\|_{C^0}, \\ &\leq C'_d \frac{1}{m_{\bar{j}}^{r_{\bar{j}}}} \|f\|_{C^{r_{\bar{j}}}}. \end{aligned}$$

For the simplicity of notations, we denote

$$p_N(x) = P_{m_1, \dots, m_d}^{[1, \dots, d]}(f)(x),$$

where $x = (x_1, \dots, x_d)$ and $N = 2m_{\bar{j}} - 1$. Moreover, we denote $k := r_{\bar{j}}$, then

$$(4.9) \quad \|f(x) - p_N(x)\|_{C^0} \leq A_{dk} N^{-k} \|f(x)\|_{C^k},$$

where A_{dk} is a constant depending on d and k .

5. C^ω DESTRUCTION OF ALL OF THE LAGRANGIAN TORI

Similar to Herman's construction, we consider C^∞ function $\tilde{T}_n(x)$ as follow:

$$\tilde{T}_n(x) = \begin{cases} \tilde{T}_n^+(x), & x \in [0, \pi]^d, \\ -\tilde{T}_n^-(x), & x \in [-\pi, 0]^d, \\ 0, & \text{others.} \end{cases}$$

$\tilde{T}_n^+(x)$ satisfies:

$$\begin{cases} \text{supp } \tilde{T}_n^+(x) \subset [0, \pi]^d, \\ \max \tilde{T}_n^+(x) = 1, \end{cases}$$

$\tilde{T}_n^-(x)$ satisfies:

$$\begin{cases} \text{supp } \tilde{T}_n^-(x) = B_{R_n}(x_0), \\ \max \tilde{T}_n^-(x) = n, \\ R_n \sim \left(\frac{1}{n}\right)^{\frac{1}{d}}, \\ x_0 = \left(-\frac{\pi}{2}, \dots, -\frac{\pi}{2}\right). \end{cases}$$

Moreover, we require $\int_{\mathbb{T}^d} \tilde{T}_n(x) dx = 0$. By Lemma 4.1, there exists a trigonometric polynomial $p_N(x_1, \dots, x_d)$ in x_l ($1 \leq l \leq d$) of degree at most N such that

$$(5.1) \quad \left\| \tilde{T}_n(x_1, \dots, x_d) - p_N(x_1, \dots, x_d) \right\|_{C^0} \leq A_{dk} N^{-k} \left\| \tilde{T}_n(x_1, \dots, x_d) \right\|_{C^k}.$$

By the construction of \tilde{T}_n , we have

$$(5.2) \quad \left\| \tilde{T}_n(x_1, \dots, x_d) \right\|_{C^k} \sim n^{\frac{k}{d}+1}.$$

Then, choosing N large enough such that

$$(5.3) \quad A_{dk} N^{-k} \left\| \tilde{T}_n(x_1, \dots, x_d) \right\|_{C^k} < \sigma \ll 1,$$

where σ is a small enough positive constant. Hence, we have

$$(5.4) \quad \begin{cases} \max p_N(x) \sim n, & \text{attained on } B_{R_n}(x_0), \\ \max p_N(x) \sim 1, & \text{on } [-\pi, \pi]^d \setminus B_{R_n}(x_0). \end{cases}$$

By (5.2) and (5.3), we have

$$(5.5) \quad N > \left(\frac{A_{dk}}{\sigma} \right)^{\frac{1}{k}} n^{\frac{1}{d} + \frac{1}{k}} \geq C n^{\frac{1}{d} + \frac{1}{k}},$$

where C is a constant independent of n . Since $\int_{\mathbb{T}^d} \tilde{T}_n(x) dx = 0$, it follows from (5.1) that $\int_{\mathbb{T}^d} p_N(x) dx = 0$. We consider the normalized trigonometric polynomial

$$(5.6) \quad \tilde{p}_N(x) = \frac{1}{n^{1-\varepsilon}} \frac{p_N(x)}{\max |p_N(x)|},$$

where $x = (x_1, \dots, x_d)$. It is easy to see that $\int_{\mathbb{T}^d} \tilde{p}_N(x) dx = 0$ and

$$(5.7) \quad \begin{cases} \max \tilde{p}_N(x) \sim \frac{1}{n^{2-\varepsilon}}, \\ \min \tilde{p}_N(x) \sim -\frac{1}{n^{1-\varepsilon}}. \end{cases}$$

It follows from (5.7) that

$$-\frac{1}{2} \min \tilde{p}_N(x) \sim \frac{1}{n^{1-\varepsilon}},$$

$$\sqrt{\max \tilde{p}_N(x)} + O(\max \tilde{p}_N(x)) \sim \frac{1}{n^{1-\frac{\varepsilon}{2}}}.$$

Hence, for n large enough, we have

$$(5.8) \quad -\frac{1}{2} \min \tilde{p}_N(x) > \sqrt{\max \tilde{p}_N(x)} + O(\max \tilde{p}_N(x)).$$

Next, we estimate $\|\tilde{p}_N(x)\|_{C^r}$. By a simple calculation, we have

$$(5.9) \quad \|p_N(x)\|_{C^r} \leq CnN^{r+1}.$$

Then,

$$\begin{aligned} \|\tilde{p}_N(x)\|_{C^r} &\leq \frac{1}{n^{1-\varepsilon}} \frac{1}{\max |p_N(x)|} \|p_N(x)\|_{C^r}, \\ &\leq \frac{1}{n^{2-\varepsilon}} \cdot CnN^{r+1}, \\ &= C \frac{1}{n^{1-\varepsilon}} N^{r+1}. \end{aligned}$$

To achieve $n^{-(1-\varepsilon)}N^{r+1} \rightarrow 0$ as $n \rightarrow \infty$, it suffices to make

$$\frac{1}{n^{1-\varepsilon}} N^{r+1} \leq \frac{1}{n^\varepsilon}.$$

Hence, we have

$$r \leq \log_N n^{1-2\varepsilon} - 1.$$

From (5.5), we take

$$N \sim n^{\frac{1}{d} + \frac{1}{k}},$$

then, we have

$$(5.10) \quad \max_N r \leq \left(\frac{1}{d} + \frac{1}{k}\right)^{-1} (1 - 2\varepsilon) - 1.$$

Since k can be made large enough, we take $k = \frac{d}{2\varepsilon}$. Let

$$\delta = \frac{4\varepsilon}{1 + 2\varepsilon} d,$$

then, we have

$$(5.11) \quad \max_N r \leq d - 1 - \delta.$$

It follows from (5.11) that $\tilde{p}_N(x) \rightarrow 0$ as $n \rightarrow \infty$ in the $C^{d-1-\delta}$ topology for any $\delta > 0$.

Since $\int_{\mathbb{T}^d} \tilde{p}_N(x) dx = 0$, then there exists the unique function $\tilde{\Psi}_n \in C^\omega(\mathbb{T}^d, \mathbb{R})$ such that

$$\int_{\mathbb{T}^d} \tilde{\Psi}_n(x) dx = 0 \quad \text{and} \quad \frac{1}{d} \Delta \tilde{\Psi}_n(x) = \tilde{p}_N(x).$$

By Schauder estimates one knows that for any $\delta > 0$, $\tilde{\Psi}_n(x) \rightarrow 0$ as $n \rightarrow \infty$ in the $C^{d+1-\delta}$ topology. According to (5.8), we have that the symplectic twist maps generated by the generating function $\tilde{h}_n(x, x') = \frac{1}{2}(x - x')^2 + \tilde{\Psi}_n(x')$ do not admit any Lagrangian tori for n large enough.

So far, we prove the following theorem

Theorem 5.1 *All of the Lagrangian tori of an integrable symplectic twist map with $d \geq 2$ degrees of freedom can be destroyed by C^ω perturbations of the generating function and the perturbations are arbitrarily small in the $C^{d+1-\delta}$ topology for a small given constant $\delta > 0$.*

Based on the correspondence between symplectic twist maps and Hamiltonian systems, together with the toy model corresponding to the case with $d = 1$, we finish the proof of Theorem 1.1.

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