Comment on "Speed Limit for Classical Stochastic Processes"

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In recent letter [Phys. Rev. Lett 121, 070601 (2018)], the speed limit for classical stochastic Markov processes is considered, and a trade-off inequality between the speed of the state transformation and the entropy production is given. In this comment, a more accurate inequality will be presented.

In recent letter of Shiraishi et al. [\[1\]](#page-0-0), for classical stochastic Markov processes, a trade-off inequality between the speed of the state transformation and the entropy production is presented. This result is very important for the study of optimization of heat engines, or even general stochastic Markov processes [\[2](#page-0-1)[–4\]](#page-0-2). In this comment, by using a new mathematical inequality, I will give an improved result.

For a stochastic Markov process with discrete states governed by the following equation,

$$
\frac{d}{dt}p_i(t) = \sum_{j(\neq i)} [W_{ij}(t)p_j(t) - W_{ji}(t)p_i(t)],
$$
\n(1)

the total variation distance between $p(0)$ and $p(\tau)$ is defined as

$$
L(\mathbf{p}(0), \mathbf{p}(\tau)) = \sum_{i} |p_i(0) - p_i(\tau)|.
$$
 (2)

Where the transition rate matrix satisfies the normalization condition $\sum_i W_{ij}(t) = 0$ and non-negativity $W_{ij}(t) \ge$ 0 for $i \neq j$. In [\[1](#page-0-0)], it is shown that

$$
\tau_I := \frac{L(\mathbf{p}(0), \mathbf{p}(\tau))^2}{2\Sigma \langle A \rangle_\tau} \leq \tau.
$$
 (3)

Where $\Sigma := \int_0^{\tau} dt \dot{\Sigma}(t)$, and $\langle A \rangle_{\tau} := \frac{1}{\tau} \int_0^{\tau} dt A(t)$ with

$$
\dot{\Sigma}(t) = \frac{1}{2} \sum_{i} \sum_{j(\neq i)} \left[(W_{ji} p_i - W_{ij} p_j) \ln \frac{W_{ji} p_i}{W_{ij} p_j} \right],
$$

$$
A(t) = \sum_{i} \sum_{j(\neq i)} [W_{ji}(t) p_i(t)].
$$
 (4)

I found that, instead of the lower bound τ_I of τ as given in [\(3\)](#page-0-3), there exists a more accurate one,

$$
\hat{\tau}_I := \frac{2L(\mathbf{p}(0), \mathbf{p}(\tau))^2}{\Sigma \langle B \rangle_\tau} \leq \tau,\tag{5}
$$

in which $\langle B \rangle_{\tau} := \frac{1}{\tau} \int_0^{\tau} dt B(t)$ with $B(t) := \sum_i B_i(t)$ and $B_i(t) := \sum_{j(\neq i)} \left[\sqrt{W_{ji}(t)p_i(t)} + \sqrt{W_{ij}(t)p_j(t)} \right]^2$. Obviously $\langle B \rangle_{\tau} \leq 4 \langle A \rangle_{\tau}$, therefore $\tau_I \leq \hat{\tau}_I \leq \tau$. The derivation of inequality [\(5\)](#page-0-4) is as follows.

$$
L(\mathbf{p}(0), \mathbf{p}(\tau)) \leq \sum_{i} \int_{0}^{\tau} dt \left| \frac{d}{dt} p_i \right|
$$

\n
$$
= \sum_{i} \int_{0}^{\tau} dt \left| \sum_{j(\neq i)} [W_{ij}(t)p_j(t) - W_{ji}(t)p_i(t)] \right|
$$

\n
$$
\leq \int_{0}^{\tau} dt \sum_{i} \sqrt{B_i(t)S_i(t)} \leq \int_{0}^{\tau} dt \sqrt{B(t)S(t)}
$$

\n
$$
\leq \left(\int_{0}^{\tau} dt B(t) \right)^{1/2} \left(\int_{0}^{\tau} dt S(t) \right)^{1/2}
$$

\n
$$
\leq \sqrt{\tau \langle B \rangle_{\tau} \Sigma / 2},
$$
 (6)

where $S_i(t) := \sum_{j(\neq i)} \left[\sqrt{W_{ij}(t)p_j(t)} - \sqrt{W_{ji}(t)p_i(t)} \right]^2$ and $S(t) := \sum_i S_i(t)$. In the third and forth lines of [\(6\)](#page-0-5), the Schwarz inequality is used, and in the last line the inequality $(a - b) \ln(a/b) \ge 4(\sqrt{a} - \sqrt{b})^2$ is used. From [\(6\)](#page-0-5), inequality [\(5\)](#page-0-4) can be obtained easily.

For the example discussed in [\[1\]](#page-0-0), transition rates $W_{10} = 1$, $W_{01} = (4\tau + 1)/(2\tau - t) - 1$, probability $p_1(t) =$ $1/2 - t/(4\tau)$, $p_0(t) = 1/2 + t/(4\tau)$. Which give that $\langle A \rangle_{\tau} =$ $5/4 + 1/(4\tau), L(\mathbf{p}(0), \mathbf{p}(\tau)) = 1/2, \Sigma = [(3\tau + 1)\ln(3\tau + 1) (2\tau+1)\ln(2\tau+1) - 3\tau\ln(3\tau) + 2\tau\ln(2\tau)/(4\tau)$, and $\langle B \rangle_{\tau} =$ $2\langle A \rangle_{\tau} + 4C$ with $C = 1/(4\tau^2) \int_0^{\tau} dt \sqrt{(t + 2\tau + 1/2)^2 - 1/4}.$ Obviously, $4C < 1/\tau^2 \int_0^{\tau} dt (t + 2\tau + 1/2) = 5/2 + 1/(2\tau) =$ $2\langle A\rangle_{\tau}$, therefore $\langle B\rangle_{\tau} < 4\langle A\rangle_{\tau}$ and consequently $\tau_I < \hat{\tau}_I$. See Fig. [1](#page-1-0) for the plots of difference $\hat{\tau}_I - \tau_I$, as well as the relative errors $(\tau - \hat{\tau}_I)/\tau$ and $(\tau - \tau_I)/\tau$.

(2018).

- [4] A. P. Solon and J. M. Horowitz, Phys. Rev. Lett. **120**, 180605 (2018).
- [1] N. Shiraishi, K. Funo, and K. Saito, Phys. Rev. Lett. **121**, 070601 (2018).
- [2] T. R. Gingrich, J. M. Horowitz, N. Perunov, and J. L. England, Phys. Rev. Lett. 116, 120601 (2016).
- [3] P. Pietzonka and U. Seifert, Phys. Rev. Lett. 120, 190602

FIG. 1: The difference $\hat{\tau}_I - \tau_I$ between the improved lower bound $\hat{\tau}_I$ obtained in this comment and the lower bound τ_I obtained in [\[1\]](#page-0-0) (left), and the relative errors $(\tau - \hat{\tau}_I)/\tau$ and $(\tau - \tau_I)/\tau$ (right).