Atomic-scale 3D structural dynamics and functional degradation of Pt alloy nanocatalysts

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Pt-based electrocatalysts are the primary choice for fuel cells due to their superior oxygen reduction reaction (ORR) activity. To enhance ORR performance and durability, extensive studies have investigated transition metal alloying, doping, and shape control to optimize the three key governing factors for ORR: geometry, local chemistry, and strain of their surface and subsurface. However, systematic optimization remains incomplete, as it requires an atomic-scale understanding of these factors and their dynamics over potential cycling, as well as their relationship to ORR activity. Here, we implement neural network-assisted atomic electron tomography to measure the 3D atomic structural dynamics and their effects on the functional degradation of PtNi alloy catalysts. Our results reveal that PtNi catalysts undergo shape changes, surface alloying, and strain relaxation during cycling, which can be effectively mitigated by Ga doping. By combining geometry, local chemistry, and strain analysis, we calculated the changes in ORR activity over thousands of cycles and observed that Ga doping leads to higher initial activity and greater stability. These findings offer a pathway to understanding 3D atomic structural dynamics and their relation to ORR activity during cycling, paving the way for the systematic design of durable, high-efficiency nanocatalysts.

Main

Introduction

Electrocatalysts are crucial substances in electrochemical energy conversion and storage, greatly enhancing the efficiency and rate of reactions in applications such as fuel cells, electrolyzers, and batteries. Especially for polymer electrolyte membrane fuel cells (PEMFCs), the effectiveness of electrocatalysts is essential to improve the slow kinetics of the oxygen reduction reaction (ORR) at the cathode, and its proper development can make it the most attractive power source for zero-emission vehicles due to their environmental benefits, energy efficiency, and power density 1,2 .

Traditionally, Pt has been commercially used as a catalyst; however, due to its high cost and limited availability, extensive research has been conducted on the development of electrocatalysts through approaches such as shape control³⁻⁵, Pt-M (M = transition metal, e.g., Ni, Co, Fe) alloys⁶⁻¹¹, dealloying^{12,13}, and doping^{9,14}. These synthesis strategies aim to enhance ORR activity while ensuring stability during fuel cell reactions. Achieving this requires fine-tuning the binding energy of intermediates (O and OH) to an optimal level at the catalyst surface⁷, which is accomplished through the appropriate interplay of geometry^{15–17}, local chemistry^{18,19}, and the strain^{5,20} of the surface and subsurface. Furthermore, from the perspective of durability, understanding the surface and subsurface structural dynamics during potential cycling and their relation to ORR activity is crucial. In this regard, precise measurement of the three-dimensional (3D) atomic structural dynamics of nanocatalysts is essential to systematically optimize catalyst performance in terms of ORR activity and durability.

Most microscopic and spectroscopic techniques provide either ensemble-averaged information over a large number of nanocatalysts²¹⁻²³ or are limited to low-dimensional measurements^{14,24-26}, instead of direct 3D local atomic structural information. Consequently, changes in the 3D atomic arrangement on actual surfaces and subsurfaces of nanocatalysts during potential cycles remain experimentally unexplored. The majority of catalytic studies rely on theoretical models and simulations^{7,27–29}, which have difficulty accurately describing the physical properties of actual structures, particularly for heterogeneous surfaces and adsorbates 30,31 .

Atomic electron tomography (AET) has recently emerged as a powerful technique for 3D structural imaging at the single-atom level^{32–35} and has been utilized to analyze Pt-based nanocatalysts^{36,37}, enabling the identification of active sites and the examination of the effects of surface strain, ligands, and doping³⁸. Additionally, neural network-assisted AET has been developed to address artifacts from data imperfections, enhancing the accuracy of atomic structures, particularly on nanocrystal surfaces $36,39$.

Here, we implemented neural network-assisted AET to determine the 3D atomic structural dynamics of octahedral PtNi nanocatalysts and clarified the effect of Ga doping in terms of geometry, local chemistry, and strain over potential cycling. By integrating these factors, we calculated how ORR activity changes over thousands of cycles.

Results

Determination of 3D atomic structural dynamics of PtNi and Ga-PtNi nanoparticles

Octahedral-shaped PtNi and Ga-doped PtNi (Ga-PtNi) nanoparticles of approximately 8-10 nm in diameter were synthesized via a one-pot synthesis method (Methods). The PtNi and Ga-PtNi particles were drop-cast onto carbon membranes with a thickness of 3-4 nm, followed by potential cycling for various numbers of cycles. A total of 8 PtNi nanoparticles (2 pristine and 6 after 12k cycles) and 9 Ga-PtNi nanoparticles (3 pristine, 2 after 4k cycles, 2 after 8k cycles, and 2 after 12k cycles) were selected for AET measurements conducted using annular dark-field scanning transmission electron microscopy (ADF-STEM) mode (see Methods and Supplementary Figs. 1-10). Note that various cycling was performed on different nanoparticles, rather than on a single nanoparticle. After image post-processing and tilt-series alignment, the atomic resolution 3D tomograms of the 8 PtNi and 9 Ga-PtNi nanoparticles were reconstructed using the GENFIRE algorithm⁴⁰. Subsequently, neural network-based volume data augmentation was applied to enhance the reliability of the atomic structures (Methods). The 3D atomic coordinates and chemical species of Pt and Ni atoms in each nanoparticle were identified from the final tomograms using atom tracing and species classification techniques (Methods)^{35,37,41}. The average precision of the 3D atomic coordinates across all nanoparticles, as determined from multislice simulations, was calculated to be 31.7 ± 5.2 pm (see Methods and Supplementary Tables 1 and 2). Note that Ga and Ni atoms cannot be distinguished in our ADF-STEM image contrast due to their similarity in atomic number.

Geometrical dynamics of PtNi and Ga-PtNi nanoparticles

Figure 1a-c depicts the 3D atomic structures and chemical compositions of six representative nanoparticles (one selected from each of the PtNi pristine, PtNi 12k, Ga-PtNi pristine, Ga-PtNi 4k, Ga-PtNi 8k, and Ga-PtNi 12k particles), clearly illustrating the heterogeneous chemical distribution within the nanoparticles. Previous studies consistently reported that undoped octahedral PtNi nanocatalysts suffer from substantial leaching of Ni atoms, resulting in a geometrical change from an octahedral to a spherical shape during the potential cycles^{26,42}. However, among the six PtNi particles measured after 12k cycles, only one particle showed the expected behavior. The catalytic activity of nanocrystals during cycling can vary significantly depending on the electrical boundary conditions connecting the catalysts to the electrode, resulting in a considerable fraction of nanoparticles being inactive during the reaction. Therefore, we selected the particle that showed the expected behavior for detailed analysis, and the results from the other particles are also summarized in Supplementary Figs. 11 and 12.

As can be seen in Fig. 1a, the PtNi nanocrystals underwent a geometrical change from an octahedral to a more spherical shape (truncated octahedron) during 12k potential cycles. Utilizing the full 3D atomic structure, we quantitatively determined the 3D surfaces of the particles via alpha-shape algorithm and calculated their sphericity (Methods). The result corroborates the qualitative observation of shape change, as described in Fig. 1d. Furthermore, for each surface atom, we applied a cotangent discretization method⁴³ to identify the facet it belongs to (Methods). Fig. 1e-g shows how the proportions of {111}, {100}, and {110} facets change over potential cycling. For PtNi, the fraction of {111} facets decreased, while those of the {100} and {110} facets increased, indicating a transformation into a truncated octahedron due to potential cycling. In contrast, for Ga-PtNi, the proportion of the {111} facet remained mostly constant during the cycle, maintaining the octahedral shape.

For PtNi catalysts, octahedral geometry is highly beneficial for catalytic activity, since {111} facets typically show higher ORR activity by a factor of 100 or more compared to other facets. Thus, this observed structural change is considered one of the main causes of degradation^{44,45}. Doping with different metallic species such as Mo or Ga has been suggested as a solution^{9,14}. Our findings for Ga-PtNi are consistent with previous reports, clearly showing that the octahedral shape is maintained throughout potential cycles up to 12k for all 9 Ga-PtNi particles measured in our study. The {111} facets of Ga-PtNi after 12k cycles exhibit more concave curvature compared to those of other Ga-PtNi particles (Fig. 1c), also consistent with a previous study²⁶, further supporting the validity of our dynamic measurements.

Additionally, the octahedral shape of nanoparticles can appear spherical in 2D projection-based imaging along certain angles (Supplementary Fig. 13), which can be misleading. This highlights the importance of conducting full 3D structural investigations to truly understand the geometric details of nanocatalysts.

Local chemistry dynamics of PtNi and Ga-PtNi nanoparticles

During potential cycling, it is known that transition metals in Pt-based alloy catalysts tend to leach out from the nanocatalyst surface, and doping with metals such as Mo or Ga can mitigate the metal dissolution^{9,14}. The averaged elemental composition fraction obtained from our 3D structures shows the expected behavior, with a substantial reduction in Ni content from 39% to 22% (a 44% reduction) for undoped PtNi during 12k cycling, while that of Ga-PtNi decreased only slightly from 42% to 35% (a 17% reduction) during the same number of cycles (see Fig. 2a). One of the main advantages of 3D structural analysis compared to 2D projection-based studies is that the physical properties can be quantitatively profiled as a function of distance from the surface, elucidating the in-depth dynamics of local chemistry. Figure 2b,c shows how the elemental composition changes as a function of distance from the surface during potential cycling for PtNi and Ga-PtNi. In their pristine states, both PtNi and Ga-PtNi exhibit a higher Ni fraction in the core regions (> 5 Å from the surface) compared to the surface/subsurface (≤ 5 Å from the surface). For undoped PtNi, the Ni fraction decreases in both the surface/subsurface and core regions, with the core region losing more Ni compared to the surface/subsurface region. In the Ga-PtNi case, the depth profile does not show a meaningful change up to 4k cycles. For cycles above 8k, a substantial reduction of Ni is observed only in the core region compared to particles that underwent fewer cycles, while the Ni fraction at the surface/subsurface is maintained. Our results confirm that Ni atoms are leaching out during cycling, and Ga doping can effectively suppress this. The notable difference between PtNi and Ga-PtNi in terms of surface Ni dynamics can be explained by the distribution of Ga dopants. The Ga atoms are expected to be mainly distributed on the surface of the nanocrystals, making it difficult for surface/subsurface Ni atoms to leach out¹⁴. Therefore, the surface/subsurface Ni fraction is mostly preserved even after 12k cycles, while core Ni atoms are leached out only through limited surface/subsurface areas where Ga atoms are absent. This preservation allows the octahedral geometry to be maintained even after extensive cycling. In contrast, the surface/subsurface Ni atoms of undoped PtNi nanocrystals, lacking the protection of Ga dopants, experience a substantial reduction in Ni fraction at both the surface/subsurface and core levels, leading to geometrical deformation towards a more spherical shape.

Note that the dissolution of core Ni atoms over potential cycling has been posited in previous experimental²⁶ and simulation studies⁴⁶. When Ni dissolution occurs on the surface/subsurface, core atoms move to fill the resulting Ni vacancies. Because Ni atoms are expected to have a lower diffusion activation barrier compared to Pt, core Ni atoms will continuously move towards the surface/subsurface and leach out. Our results directly show this phenomenon at the atomic scale in 3D.

The migration of Ni atoms during cycling not only affects the geometrical deformation of the nanocrystals but also directly influences the ORR activity via the ligand effect⁴⁷. Therefore, based on our 3D atomic structures, we calculated the chemical short-range order parameter (CSRO) for every atom within the nanoparticles to describe their local chemical disorder (Methods). A CSRO value close to 0 indicates that the local region around a given atom exhibits the chemical composition of the overall nanoparticle. Values greater than 0 suggest local chemical segregation, and values less than 0 indicate a degree of alloying higher than the global composition. Undoped PtNi exhibits Pt CSRO values of about 0.3 at the surface, indicating a strong tendency for Pt segregation (Fig. 2d). Moving towards the interior of the particle, the Pt CSRO approaches 0, suggesting that the local composition becomes similar to the global composition. After 12k cycles, the Pt CSRO at the surface decreases to less than half of the value of the pristine particle (indicating alloying), while the Pt CSRO of the core region becomes substantially higher than before cycling (indicating segregation).

Pristine Ga-PtNi exhibits behavior similar to undoped pristine PtNi, with surface/subsurface Pt segregation and a more alloyed core (Fig. 2e). As the number of cycles increases, the surface/subsurface Pt CSRO gradually decreases (indicating more alloyed configurations), while the core Pt CSRO gradually increases (indicating more segregated configurations). The difference in the Pt CSRO depth profile between pristine and 12k cycled Ga-PtNi particles is less pronounced than in undoped PtNi. After 12k cycles, the surface/subsurface and core Pt CSRO values of Ga-PtNi are at the same level, whereas undoped PtNi shows an inverted behavior in Pt CSRO for the surface/subsurface and core regions between pristine and 12k cycled particles.

We expect that the CSRO of Ni will behave approximately opposite to the Pt CSRO, as shown in Fig. 2f,g. In the pristine state, Ni exists in a more alloy-like form compared to Pt at the surface/subsurface. However, core Ni tends to be in a more segregated state. As cycling progresses, this trend changes, with core Ni showing an alloy composition consistent with the entire nanoparticle. This indicates that segregated Ni atoms are diffusing towards the surface during cycling. Note that Ni CSRO and Pt CSRO are not completely complementary; the CSRO for Pt, which has a higher global composition, is more sensitive to local chemistry, as shown in Fig 2.d-g.

Strain dynamics of PtNi and Ga-PtNi nanoparticles

Strain is another crucial factor in determining the performance of a catalyst. Strain alters the electronic structure, resulting in changes in adsorption energies. Pt nanoparticles are expected to show their optimal ORR activity at approximately 3% compressive surface strain^{14,48}. We calculated the volumetric strains of Pt atoms by comparing the measured local lattice constant with that of bulk Pt^{49} (Methods). As shown in Fig. 3a, both PtNi and Ga-PtNi exhibit compressive surface strain in their pristine state, with Ga-doped particles showing stronger compressive strain closer to the target value of 3%, indicating that Ga doping is beneficial for improving the catalytic performance.

As the cycling continues, the strain experiences a substantial change. As shown in Fig. 3a, after 12k cycles, strain relaxations of approximately 40.9% (from −2.2% to −1.3%) and 26.1% (from −2.3% to −1.7%) were observed for PtNi and Ga-PtNi, respectively. This result is consistent with the findings from the local chemistry analysis: Since Pt has a larger lattice constant than $Ni⁴⁹$, the dissolution of surface/subsurface and core Ni over the cycles results in strain relaxation toward a less compressive configuration. Again, since Ga doping effectively prevents the dissolution of Ni atoms, the strain relaxation for Ga-PtNi is also suppressed compared to the undoped particles.

We also quantitatively profiled the strain as a function of distance from the surface, as illustrated in Fig. 3b,c. For undoped PtNi, strain relaxation occurs in both the surface/subsurface and core regions during potential cycles, indicating that the Ni dissolution occurs throughout the nanoparticle (Fig. 3b). However, for Ga-PtNi, strain relaxation is observed only in the core region with minimal changes in the surface/subsurface (Fig. 3c). This suppressed strain relaxation indicates that Ga doping mitigates Ni dissolution from the surface, aligning with the observed local chemistry dynamics (Fig. 2).

Calculation of the ORR activity considering geometry, local chemistry, and strain altogether

Finally, we relate the geometry, local chemistry, and strain to calculate the local ORR activity $[\ln(j/j_{Pt(111)})]$ of all surface Pt atoms (Methods). By simultaneously considering geometry, local chemistry, and strain through generalized coordination number $(GCN)^{50}$, we derive a modified alloysensitive GCN for each surface Pt atom as a descriptor for calculating ORR activity from our 3D atomic structures (see Fig. 4a-f and Methods). As shown in Fig. 4a-f, a clear inhomogeneity can be observed from the ORR activity distribution. Moreover, the {111} facets exhibit better ORR activity compared to those of the $\{100\}$ or $\{110\}$ facets, consistent with previous studies⁶.

By examining the average values of each nanoparticle's $j/j_{Pt(111)}$ over cycles, we clarified the effect of potential cycling on the overall ORR behavior for PtNi and Ga-PtNi (Fig. 4g). As the cycles progress, the ORR activity decreases for both PtNi and Ga-PtNi. Specifically, the average value of $j/j_{Pt(111)}$ for PtNi shows a decrease of about 17%, which can be attributed to a combinatory effect arising from the decreased fraction of {111} facets, loss of surface Ni atoms, and strain relaxation. Whereas, the ORR activity of Ga-PtNi is substantially more robust against potential cycling, with about 4% reduction of ORR even after 12k cycles. As discussed in previous sections, Ga doping effectively suppresses changes in all three factors—geometry, chemistry, and strain—during cycling. Our descriptor accounts for these factors, successfully predicting the expected ORR behavior.

Conclusion

In this study, we observed the changes in ORR activity of PtNi and Ga-PtNi nanoparticles throughout potential cycling by combining three crucial factors for catalytic activity: geometry, local chemistry, and strain. To achieve this, we utilized neural network-assisted AET to determine the full 3D atomic structures of several nanoparticles at various stages of cycling.

Geometric analysis showed that PtNi catalysts experience unfavorable shape changes from octahedral to spherical due to cycling, which can be mitigated by Ga doping. Analysis of the local chemical composition indicated that PtNi catalysts undergo surface alloying which negatively impact ORR activity, and Ga doping prevents this effect. In terms of strain, PtNi catalysts shift towards less compressive surface strain, which is detrimental to ORR activity, whereas Ga doping helps maintain the favorable compressive strain throughout cycling. By integrating these three factors—geometry, local chemistry, and strain—we determined the ORR activity of each nanoparticle at the single-atom level, enabling us to identify the average catalytic activity of each nanoparticle and observe the distinct degradation of this activity over multiple cycles without Ga doping.

Our approach will enable the efficient design of high-performance nanocatalysts through a fundamental understanding of the relationship between surface structure and properties at the atomic scale.

Figure 1 | Experimentally determined 3D atomic structures of representative PtNi and Ga-PtNi nanoparticles during potential cycles and geometric characterization. a, Overall 3D atomic structure of different PtNi (pristine and 12k) and Ga-PtNi (pristine, 4k, 8k, and 12k) nanoparticles after the indicated number of potential cycles. **b**, 3D atomic structures after a 180-degree rotation along the [001] axis from those in (**a**), with one octant of the nanoparticles removed to reveal the internal atomic structure. **c**, 1 Å thick internal slices at the center of the 3D tomograms perpendicular to the [001] direction. The intensity is shown in grayscale, with the atomic coordinates of Pt and Ni marked by blue and red dots, respectively. A concave (111) facet (marked with white lines) is observed in the Ga-PtNi nanoparticle after 12k cycles. Scale bar, 2 nm. Note that, the representative nanoparticles shown in (**a-c**) were PtNi-pristine-p2, PtNi-12k-p1, Ga-PtNi-pristine-p1, Ga-PtNi-4k-p1, Ga-PtNi-8k-p2, and Ga-PtNi-12k-p2 (Supplementary Tables 1 and 2). **d**, Sphericity of the PtNi and Ga-PtNi nanoparticles during the potential cycles. The sphericity values of the ideal sphere, truncated octahedron, and octahedron are marked by dotted lines. **e-g**, Fraction of {111} (**e**), {100} (**f**), and {110} (**g**) facets during potential cycles. Note that the values and error bars for each cycle corresponding to (**d-g**) represent the averages and standard deviations obtained from 2 PtNi pristine, 3 Ga-PtNi pristine, 2 Ga-PtNi 4k, 2 Ga-PtNi 8k, and 2 Ga-PtNi 12k samples. We have only one particle for PtNi 12k, and the error bar cannot be defined in the same manner in this case.

Figure 2 | Quantitative characterization of the chemical compositions of PtNi and Ga-PtNi nanoparticles during potential cycles. a, Ni fraction of PtNi and Ga-PtNi nanoparticles during potential cycles. **b**, **c**, Characterization of the Ni fraction as a function of distance from the surface during potential cycles for PtNi (**b**) and Ga-PtNi (c) nanoparticles. **d**, **e**, Chemical short-range order parameter (CSRO) of Pt $(\alpha_{\text{PM}i})$ as a function of distance from the surface during potential cycles for PtNi (**d**) and Ga-PtNi (**e**) nanoparticles. **f**, **g**, CSRO of Ni (α_{NiPt}) as a function of distance from the surface during potential cycles for PtNi (**f**) and Ga-PtNi (**g**) nanoparticles. A CSRO value close to 0 indicates that the local region around the atom reflects the overall nanoparticle's chemical composition (represented by a dotted line), while values greater than 0 suggest local chemical segregation and values less than 0 indicate a higher degree of alloying than the global composition. Note that the values and error bars for each cycle corresponding to (**a-g**) represent the averages and standard deviations obtained from 2 PtNi pristine, 3 Ga-PtNi pristine, 2 Ga-PtNi 4k, 2 Ga-PtNi 8k, and 2 Ga-PtNi 12k samples. We have only one particle for PtNi 12k, and the error bar cannot be defined in the same manner in this case.

Figure 3 | Quantitative characterization of the strain of PtNi and Ga-PtNi nanoparticles during potential cycles. a, Local volumetric strain of surface Pt atoms during potential cycles of PtNi and Ga-PtNi nanoparticles. **b**, **c**, Strain of Pt as a function of distance from the surface during potential cycles for PtNi (**b**) and Ga-PtNi (**c**) nanoparticles. Note that the values and error bars for each cycle corresponding to (**a-c**) represent the averages and standard deviations obtained from 2 PtNi pristine, 3 Ga-PtNi pristine, 2 Ga-PtNi 4k, 2 Ga-PtNi 8k, and 2 Ga-PtNi 12k samples. We have only one particle for PtNi 12k, and the error bar cannot be defined in the same manner in this case.

Figure 4 | Quantitative characterization of ORR activity of PtNi and Ga-PtNi nanoparticles during potential cycles. a-f, The 3D ORR activity $[\ln(j/j_{Pt(111)})]$ distribution of PtNi pristine (**a**), PtNi 12k (**b**), Ga-PtNi pristine (**c**), Ga-PtNi 4k (**d**), Ga-PtNi 8k (**e**), and Ga-PtNi 12k (**f**). Note that, the representative nanoparticles shown in (**a-f**) were PtNi-pristine-p2, PtNi-12k-p1, Ga-PtNi-pristine-p1, Ga-PtNi-4k-p1, Ga-PtNi-8k-p2, and Ga-PtNi-12k-p2 (Supplementary Tables 1 and 2). **g**, Average of $(j/j_{Pt(111)})$ during potential cycles of PtNi and Ga-PtNi nanoparticles. Note that the values and error bars for each cycle corresponding to (**g**) represent the averages and standard deviations obtained from 2 PtNi pristine, 3 Ga-PtNi pristine, 2 Ga-PtNi 4k, 2 Ga-PtNi 8k, and 2 Ga-PtNi 12k samples. We have only one particle for PtNi 12k, and the error bar cannot be defined in the same manner in this case.

METHODS

Sample preparation

Materials

Platinum(II) acetylacetonate (Pt(C₅H₇O₂)₂, 97%), nickel(II) acetylacetonate (Ni(C₅H₇O₂)₂, 95%), and gallium(III) acetylacetonate (Ga(C₅H₇O₂)₃, 99.99%) were purchased from Sigma Aldrich. N,Ndimethylformamide $(C_3H_7NO, 99.5\%)$ was obtained from Junsei Chemical.

Synthesis of PtNi

32 mg of platinum(Ⅱ) acetylacetonate and 16 mg of nickel(Ⅱ) acetylacetonate were dissolved in N,Ndimethylformamide with vigorous stirring for 1 hour. The mixture was then heated to 130 ℃, maintained at that temperature for 42 hours, and naturally cooled to room temperature. After natural cooling, the precipitate was centrifuged with a mixture of ethanol and acetone and then dried at 60 ℃ in a vacuum overnight.

Synthesis of Ga-PtNi

The synthesis procedure was the same as that of PtNi noted above. Subsequently, 6 mg of gallium(Ⅲ) acetylacetonate was added to the mixture, which was reheated to 130 ℃ and maintained for another 42 hours. After natural cooling, the precipitate was centrifuged with a mixture of ethanol and acetone and then dried at 60 ℃ in a vacuum overnight. These synthesis conditions are known to result in approximately 1.5% Ga doping¹⁴.

Electrochemical potential cycling

The resulting PtNi and Ga-PtNi nanoparticles were dispersed in ethanol with 3 hours of sonication. The dispersions were then drop-cast onto carbon membranes with a thickness of 3-4 nm and dried in air at room temperature. To perform potential cycling for various numbers of cycles, at least one TEM grid was prepared for each cycle and doping configuration (pristine and 12k cycles for undoped PtNi, and pristine, 4k, 8k, and 12k cycles for Ga-PtNi).

For potential cycling, each TEM grid was directly loaded onto a glassy carbon electrode (GCE) using a Teflon cap. The TEM grid on the GCE was subjected to potential cycles in the range from 0.6 to 1.0 V_{RHE} in an oxygen (O₂)-saturated 0.1 M HClO₄ solution at a scan rate of 100 mV s⁻¹, as an accelerated stress test. Potential cycling was conducted in a three-electrode system using a potentiostat (PGSTAT302N, Metrohm Autolab). A platinum wire and a saturated calomel electrode (Hg/Hg_2Cl_2) were employed as the counter electrode and a reference electrode, respectively. All potentials were calibrated to the reversible hydrogen electrode (RHE) in a hydrogen (H₂)-saturated 0.1 M HClO₄. After the potential cycles, the samples were annealed in a vacuum at 150 °C for 10 hours to reduce hydrocarbon contamination.

STEM data acquisition

The tomographic tilt series were acquired using three different microscopes: an FEI Titan Themis Z at the Institute of Next-generation Semiconductor convergence Technology (INST), the TEAM 0.5 microscope equipped with the TEAM stage at the National Center for Electron Microscopy (NCEM), and an FEI Spectra Ultra TEM at the KAIST Analysis Center for Research Advancement (KARA). From the tilt series datasets, atomic resolution electron tomograms were obtained for nanoparticles of 2 PtNi pristine, 6 PtNi 12k, 3 Ga-PtNi pristine, 2 Ga-PtNi 4k, 2 Ga-PtNi 8k, and 2 Ga-PtNi 12k cycles (see Supplementary Tables 1 and 2). Note that various cycling was performed on different nanoparticles, not on the same nanoparticle. The images for each tilt series were obtained using the ADF-STEM mode with 200 kV or 300 kV acceleration voltage. To correct the effect of stage drift during measurement, three or four consecutive images of 1024×1024 pixels were collected using a dwell time of 3 μs at each tilt angle (see Supplementary Tables 1 and 2 for detailed microscope parameters). Each tilt series was measured using a total electron dose in the range of 2.1 × 10⁵ *e* Å⁻² to 4.0 × 10⁵ *e* Å⁻². To ensure minimal structural changes induced by the electron beam, zero-degree projections were measured three times: at the beginning, in the middle, and at the end of the tilt series acquisition for comparison (see Supplementary Figs. 9 and 10). As shown in Supplementary Figs. 9 and 10, no substantial structural changes were observed, indicating that the PtNi nanoparticles remain stable within the given total electron dose range.

Image post-processing

We performed image post-processing for the tilt series using a series of procedures, including drift correction, scan distortion correction, image denoising⁵¹, background subtraction, and tilt-series alignment based on center-of-mass and common-line alignment, following the methods described in previous works^{34-36,41,52-54}.

(I) Drift and scan distortion correction: We performed drift correction on the images at each tilt angle by estimating the linear stage drift from three or four consecutively acquired ADF-STEM images and compensating for it using an affine transform. Subsequently, the scan distortion of the images was corrected using a scan distortion matrix estimated from images of a single-crystal silicon (110) standard sample. The three or four consecutive images, after drift and scan distortion correction, were then averaged to create an experimental image for each tilt angle.

(II) Image denoising: To remove the Gaussian-Poisson mixed noise in the ADF-STEM images, we applied the block-matching and 3D filtering (BM3D) algorithm⁵¹ using Gaussian-Poisson noise parameters estimated from the experimental tilt series.

(III) Tilt-series alignment: To remove the background signal within the nanoparticles, we defined a 2D mask slightly larger than the boundary of the nanoparticles for each projection image. We then solved the Dirichlet boundary value problem of the discrete Laplace's equation to determine the background signal, which was subsequently subtracted from each denoised image. Afterward, each tilt series was aligned with sub-pixel accuracy using the center-of-mass³³ and common-line alignment³⁵ methods.

Generation of training/test datasets for the neural network-assisted AET

Input and target datasets were prepared for training, validation, and testing of the neural networks following the previously established procedures^{36,39}. A total of 4,800 simulated tomograms were generated, with 4,000 for training, 400 for validation, and 400 for testing. The generation of input data followed a four-step process.

(1) We first created a blank 3D volume with a size of $384 \times 384 \times 384$ voxels and a voxel size of 0.329 Å. Within this blank volume, we constructed a random 3D shape with a volume ranging from 865,700 to 1,121,000 \AA^3 , and placed a randomly oriented fcc atomic structure of a 3.84 Å lattice constant inside the random 3D shape. The atomic structures included point defects, with percentages randomly chosen between 0% and 0.5%, and random spatial displacements of about 20 pm root-mean-square deviation (RMSD). The generated atomic positions were then randomly assigned as either Ni or Pt atoms, maintaining a Ni-to-Pt ratio of 1:1.

(2) The sharp 3D atomic potentials at the atomic positions were calculated^{36,39,55}, and the resulting sharp potentials were convolved with Gaussian kernels to simulate broadening effects, including thermal vibrations. The standard deviations (σ) of the Gaussian kernels for Ni and Pt atoms were randomly chosen from a Gaussian distribution with a mean of 0.60 Å for Ni and 0.70 Å for Pt, and standard deviation of 0.11 Å for both. To account for the intensity ratio of Pt and Ni atoms in experimentally reconstructed tomograms, a weighting factor of 0.7 was applied to the Pt atomic potential.

(3) 41 forward-projections of the calculated 3D volumes for tilt angles ranging from −65° to +65° were generated as a tilt series for each volume. The projection size was 384×384 pixels with a pixel size of 0.329 Å. To mimic the real experiment conditions, Poisson noise and random tilt angle errors up to $\pm 0.3^{\circ}$ were added.

(4) 3D tomograms were reconstructed from the generated tilt series and the tilt angles using the GENFIRE algorithm⁴⁰. The GENFIRE parameters of the fast Fourier transform interpolation method with 100 iterations, an oversampling ratio of 2, and an interpolation radius of 0.3 pixels were used for the reconstructions. To normalize the intensity of the 3D volumes, we first calculated the average intensity of the $3 \times 3 \times 3$ voxels at all the Pt atomic positions. Then, we divided the reconstructed volume by this averaged intensity.

For generating target data (ground truth), we followed steps (1) and (2) of the input data generation procedure. The only difference was that Gaussian kernels with the standard deviations (σ) of 0.60 Å for Ni and 0.70 Å for Pt were applied to the calculated sharp potentials. For intensity normalization, we also divided the generated target volume by the average intensity calculated in step (4).

The framework of deep learning neural network and training

The deep learning neural network used in this study follows the framework suggested in our previous work³⁹. The primary difference is the increased depth; we utilized a neural network that is two layers deeper than that of the previous study³⁹ (see Supplementary Fig. 14). As shown in Supplementary Fig. 14, both the input and output data sizes were $384 \times 384 \times 384$ voxels. The main building blocks of the network include $3 \times 3 \times 3$ convolutions with a stride of 2 for down-sampling, $2 \times 2 \times 2$ max-pooling for down-sampling, $3 \times 3 \times 3$ transposed convolutions with a stride of 2 for up-sampling, and two 3 \times 3×3 convolutions with a stride of 1. To prevent overfitting, we employed the dropout method⁵⁶. Two activation functions were used: the Leaky Rectified Linear Unit (Leaky ReLU) 57 with a leakage coefficient of 0.2 for all layers except the final output layer, and the Rectified Linear Unit (ReLU)⁵⁸ for the final output layer. The mean squared error was used as the loss function, and the Adam optimizer⁵⁹ was employed with a learning rate of 1×10^{-3} . Supplementary Figure 15 displays the learning curve for the training and validation sets, showing no signs of overfitting or divergence throughout the training process. The trained neural network model at the 30th epoch was used for the augmentation of experimental tomograms.

3D reconstruction and neural network-based volume data augmentation

After image post-processing, 3D tomograms were reconstructed from the post-processed tilt series using the GENFIRE algorithm⁴⁰. To enhance the quality of the tomograms, we applied angular refinement⁴⁰ implemented in the GENFIRE algorithm, as well as in-plane rotational and translational re-alignment³⁶. After the corrections, we reconstructed the final 3D tomograms using the following GENFIRE parameters: discrete Fourier transform interpolation method, number of iterations of 1000, an oversampling ratio of 4, and an interpolation radius of 0.1 pixels (Supplementary Table 1 and 2). However, the tomograms suffered from artifacts caused by data imperfections due to sparse sampling and the missing wedge problem. To address this, we augmented the tomograms using the deep learningbased neural network as described above^{36,39}. To normalize experimental tomograms used for neural network training, we calculated the average integrated intensity, I_{ave} , of the approximately $1.0 \times 1.0 \times$ 1.0 Å³ ($3 \times 3 \times 3$ voxels) volumes around the Pt atom positions in the experimental tomograms. Several normalization factors were tested, and dividing the raw experimental tomograms by $\frac{I_{ave}}{1.3}$ demonstrated the best performance for neural network-based volume augmentation, specifically in terms of atom tracing as discussed in the next section. Therefore, this normalization factor was consistently applied to all experimental tomograms before applying the data augmentation.

Atom tracing and species classification

For 3D tomograms before and after neural network-based volume data augmentation, we determined the 3D atomic coordinates using the following atom tracing and species classification processes $35,37,41$. (i) The 3D local maxima positions in the tomograms were identified and ranked by peak intensity in descending order. Starting with the highest intensity, a volume of $5 \times 5 \times 5$ voxels $(1.6 \times 1.6 \times 1.6$ to $1.8 \times 1.8 \times 1.8$ Å³, depending on voxel size) centered on each local maximum was extracted, and a 3D Gaussian function was fitted to the cropped volume to determine the precise peak position. The identified peak position was added to a list of potential atom positions only if it satisfied a minimum distance of 2.0 Å from any previously listed fitted positions. This procedure was repeated for all 3D local maxima, resulting in a list of potential atom positions.

(ii) A *k*-means clustering algorithm^{35,37,41} was used to classify the potential atom positions into three types of atoms (non-atom, Ni, and Pt). Atoms classified as Ni or Pt located outside the 3D boundary of the nanoparticle, as defined by the Otsu threshold, were further reclassified as non-atom. Note that due to the similarity in atomic numbers between Ga and Ni, they cannot be distinguished in AET, resulting in Ga being counted as Ni; however, since the compositional change of Ni in Ga-PtNi during cycling varies by more than 7 percentage points (from 42% to 35%), the 1.5% Ga doping does not interfere with the calculation of Ni fractions much.

(iii) To fully determine the atomic structures that were not identified in the previous steps, we applied an additional atom tracing process^{37,41}. Since the PtNi nanoparticles have an fcc crystal structure, the tomogram was first rotated to align with the crystallographic direction. The rotated 3D tomogram was sliced along the [001] direction for each atomic layer, with the slicing positions corresponding to the peak positions in the histogram of the atom position components along the [001] direction. We identified the positions of 2D local maxima, along with their corresponding intensities, for all slices. Starting with the highest intensity, the 3D Gaussian fitting procedure was performed using the 2D local maxima positions as initial estimates. The size of the cropped volume for the 3D Gaussian fitting was adjusted by varying the side length from 3 to 7 pixels. The best-fit position, determined from the volume size yielding the smallest mean squared residuals, was added to the traced atom list, provided that it satisfied a minimum distance of 2 Å from neighboring atoms. To finalize the 3D atomic structures and their chemical compositions, we repeated the process described in (ii), and as a result, obtained the full 3D atomic structures of 8 PtNi and 9 Ga-PtNi nanoparticles.

(iv) To verify the consistency between the determined atomic structures and the measured projections, we calculated the R-factors by comparing the experimentally obtained tilt series with the simulated tilt series generated from the forward projections of the final 3D atomic models. The average R-factor for all nanoparticles was determined to be 0.09 ± 0.01 (Supplementary Tables 1 and 2), a value considered acceptable in the crystallography community, indicating consistency between the two sets of images.

Assignment of experimental 3D atomic positions to ideal fcc lattices

The 3D atomic positions were assigned to ideal fcc lattice sites through the following procedures^{39,37}.

(a) A Pt atom nearest to the average position of all the 3D atomic coordinates was selected as the origin of an ideal fcc lattice.

(b) Subsequently, the nearest fcc sites to this origin atom position were calculated using the initial fcc lattice vectors, with a lattice constant of 3.9 Å. For each calculated fcc site, if an atom was located within 25% of the nearest neighbor distance of the fcc lattice, the atom was assigned to that corresponding fcc lattice site.

(c) The nearest neighbor search was repeated for all newly assigned fcc lattice sites, and this process was continued until no additional atoms could be assigned to the lattice.

(d) New fcc lattice vectors were fitted to the assigned atoms by changing the translation, 3D rotation, and lattice constant parameters, aiming to minimize the discrepancy between the 3D atomic positions and their corresponding lattice sites in the fitted fcc lattice.

(e) The processes from (a) to (d) were performed iteratively until the fitted lattice vectors remained unchanged.

Precision estimation using STEM multislice simulation

A precision estimation was conducted to evaluate the reliability of the 3D atomic models of PtNi and Ga-PtNi nanoparticles. Projection images were generated from the finalized 3D atomic models at experimental tilt angles using an efficient multislice simulation, specifically the PRISM simulation $60,61$, with an interpolation factor of 2, a slice thickness of 2 Å, sixteen frozen phonon configurations, and aberration values of -352.8 nm C₃ and 1.5 mm C₅ for nanoparticles imaged at INST, 0 nm C₃ and 5 mm C_5 for nanoparticles imaged at NCEM, and −521 nm C_3 and 0.58 mm C_5 for nanoparticles imaged at KARA. Other microscope parameters, such as electron acceleration voltages, convergence semiangles, and detector inner and outer semi-angles were adjusted to match the experimental parameters provided in Supplementary Tables 1 and 2. To account for the influence of electron probe size and other incoherent effects, a Gaussian kernel with an optimized standard deviation was applied to each PRISMsimulated image (see Supplementary Tables 1 and 2).

We applied the GENFIRE algorithm⁴⁰ to reconstruct the 3D tomograms from the PRISM-simulated tilt series, using the same parameters described above. The 3D atomic structures of Pt and Ni atoms were subsequently determined from these tomograms using previously described methods (see above for more details on atom tracing and species classification).

We estimated a constant background value of experimental tomograms by averaging the intensity of approximately $1.0 \times 1.0 \times 1.0$ Å³ ($3 \times 3 \times 3$ voxels) volumes surrounding the Pt and Ni atom positions in the experimental tomograms, and this value was added to the 3D tomogram obtained from the PRISM-simulated tilt series.

Subsequently, the neural network-based volume data augmentation process was performed as previously described, and the 3D atomic structures were finally determined. To compare the atomic structures obtained from experiments with those from PRISM simulations, we calculated the distances between atoms common to both the experimental and PRISM-simulated structures. Atom pairs with distances less than a specified threshold (half of the first nearest-neighbor distance in the ideal fitted fcc lattice, 1.4 Å) were classified as common atom pairs. The analysis revealed that, on average, 94.2 \pm 2.6% of atoms were successfully identified in the simulation. Additionally, the average root-meansquare deviation (RMSD) of all common atom pairs across all nanoparticles (i.e., the precision of the atomic coordinates; see Refs.^{34,35,52}) was determined to be 31.7 ± 5.2 pm (see Supplementary Tables 1) and 2).

Calculation of sphericity

To determine how closely each given nanoparticle's shape resembles that of a perfect sphere, we calculated the sphericity⁶², defined as $\pi^{1/3}(6V)^{2/3}/A$, where A and V represent the surface area and the volume of the nanoparticle, respectively. The surface area and volume of the nanoparticle were calculated using the alpha-shape algorithm 63 , with the lattice constant obtained from global fcc fitting used as the alpha radius.

Chemical short-range order parameter (CSRO)

To understand the local chemistry of each atom, we calculated the CSRO, which is a pairwise multicomponent short-range order parameter $64,65$. The CSRO between Pt atoms and their nearest neighbor Ni atoms (α_{PtNi}) and the CSRO between Ni atoms and their nearest neighbor Pt atoms (α_{NiPt}) are given by

$$
\begin{cases} \alpha_{\text{PtNi}} = 1 - \frac{p_{\text{PtNi}}}{c_{\text{Ni}}} \\ \alpha_{\text{NiPt}} = 1 - \frac{p_{\text{NiPt}}}{c_{\text{Pt}}} \end{cases}
$$
 (1)

where p_{PKNi} is the Ni atom fraction among the 1st neareast neighbors of Pt atoms, C_{Ni} is the Ni atom fraction in the entire nanoparticle, p_{NIPt} is the Pt atom fraction among the 1st nearest neighbors of Ni atoms, *C*Pt is the Pt atom fraction in the entire nanoparticle. A CSRO value near 0 indicates that the local region around a given atom reflects the chemical composition of the entire nanoparticle. Values above 0 suggest local chemical segregation, while values below 0 denote a higher degree of alloying than the global composition of the nanoparticle. Note that the $1st$ nearest neighbor is defined as a set of atoms within a cutoff distance determined as the midpoint between the $1st$ and $2nd$ nearest neighbors, based on the lattice constant obtained from the global fcc fitting (see 'Assignment of experimental 3D atomic positions to ideal fcc lattices' section of Methods).

Defining the surface atoms and calculating the distance from the surface

Surface atoms were defined as the atoms located outside a 3D mask that was 6 Å smaller than the nanoparticle boundary determined during the non-atom classification step described above. The distance from the surface for each atom was defined as the shortest distance to the nearest surface atom. Note that the criterion of 6 Å was determined to ensure that the surface atoms uniformly cover the entire nanoparticle.

Calculation of local volumetric strain

To obtain local strain, we performed local fcc lattice fitting using the following procedure. First, we identified a set of nearest neighbor atoms centered around each atom based on the globally fitted fcc lattice constant. Using the translation, 3D rotation, and lattice constant parameters of the globally fitted lattice as the initial condition, we performed a similar fcc lattice fitting for this set of atoms by minimizing the error between the measured atomic positions and the corresponding lattice sites. Through these determined local fcc lattice vectors, we calculated the lattice constant for each atom, which we defined as the local lattice constant (a_{local}) . Using the local lattice constant, the local volumetric strain (*S*) for each atom was calculated using the following equation:

$$
S = \left(\frac{a_{\text{local}}}{a_{\text{ref}}} - 1\right). \tag{2}
$$

Here, a_{ref} and a_{local} represent the lattice constant of bulk Pt^{49} (3.912 Å) and the local lattice constant of each atom, respectively. The positive and negative signs of the local volumetric strain indicate tensile and compressive strain, respectively.

Assignment of the facet indices

To assign facet indices for each surface atom, we first calculated the normal vector using the discrete Laplace-Beltrami operator via cotangent discretization. For each surface atom, we formed a set of atoms within a cutoff distance defined as $\frac{1}{2} \left(\frac{a_{\text{global}}}{\sqrt{2}} \right)$ $\frac{1}{\sqrt{2}} + a_{\text{global}}$, where a_{global} is the lattice constant obtained from the global fcc fitting. We then computed the convex hull of this atom set and calculated the mean curvature normal operator using the cotangent discretization method⁴³. The mean curvature normal operator, evaluated at the given surface atom position, was normalized and defined as the normal vector of the atom. We further applied 3D Gaussian kernel averaging to the normal vectors with a standard deviation of a_{global} , equivalent to the 2^{nd} nearest neighbor distance. For the kernel averaged normal vector for each surface atom, we calculated dot products with the normal vectors of three crystallographic facet families ({100}, {110}, and {111}) and assigned it to the facet with the largest dot product value.

Calculation of the ORR activity

To calculate the ORR activity for alloy nanocatalysts, we combined the strain effect and ligand effect with the GCN.

Typically, the GCN^{17} for atom *i* can be expressed as

$$
GCN(i) = \sum_{j=1}^{n_i} \frac{CN(j)}{CN_{\text{max}}}.
$$
 (3)

Here, n_i represents the coordination number for atom *i*, CN (i) denotes the coordination number of the 1st nearest neighbor atom *j* of atom *i*, and CN_{max} is the number of 1st nearest neighbors in the bulk (12 for fcc).

By incorporating a prefactor that accounts for the strain effect⁴⁸ and a correction factor for Ni (1.59) accounting for the ligand effect (estimated from density functional theory calculation⁵⁰), the extended generalized coordination number (so-called alloy-sensitive-GCN; ASGCN⁵⁰) for atom *i* can be expressed as

$$
ASGCN(i) = \frac{1}{1+S(i)} (1.59 \times GCN(i))_{Ni} + GCN(i))_{Pt}),
$$
 (4)

where $GCN(i)_{Ni}$ and $GCN(i)_{Pt}$ represent the contributions of the Ni and Pt atoms of the 1st nearest neighbors to the GCN value, respectively, therefore $(GCN(i)_{Ni} + GCN(i)_{Pt}) = GCN(i)$. Since alloysensitive GCN exhibits a volcano relationship with ORR activity $(\ln(j/j_{Pt(111)}))^{50}$, we can calculate ORR activity at the single-atom level by determining the ASGCN for each Pt surface atom.

However, in the case of the ASGCN, simulations do not accurately represent actual experiments. Specifically, when estimating the ligand effect of Ni, simulations assume that the surface atoms are composed solely of Pt and the Ni-Pt alloy can form only in the subsurface. In contrast, actual nanocatalysts do not have a perfectly homogeneous Pt surface; instead, they exhibit local chemical heterogeneity.

Therefore, the ligand effects of subsurface and surface Ni should be considered separately. In general, many simulations account for the ligand effect of subsurface Ni^{50,66}, leading us to assume that the ligand effect of surface Ni can be considered negligible compared to that of subsurface Ni. Consequently, the alloy-sensitive GCN can be revised to the modified alloy-sensitive GCN (MASGCN), as shown in Eq. (5) :

$$
MASGCN(i) = \frac{1}{1+S(i)} \left(1.59 \times GCN(i)^{Ni(subsurface)} + GCN(i)^{Ni(surface)} + GCN(i)^{Pt} \right).
$$
 (5)

The ORR calculation results in Fig. 4 are calculated based on the MASGCN above.

Data availability

All of our experimental data, tomographic reconstructions, determined atomic structures, and ORR results will be posted on a public website upon publication.

Code availability

Source codes will be posted on a public website upon publication.

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Author Contributions

Y.Y. conceived the idea and directed the study. S.L., K.L., and E.C. synthesized the PtNi/Ga-PtNi nanoparticles and conducted the potential cycling. C.J., J.L., H.J., P.E., and Y.Y. designed and performed the electron tomography experiments. J.L. and Y.Y. designed and trained the neural network. C.J., C.O., H.J., and Y.Y. conducted the data analysis. C.J., J.L., H.J., and Y.Y. wrote the manuscript. All authors commented on the manuscript.

Competing interests

C.J., J.L., H.J., S.L., K.L., E.C., and Y.Y. have patent application (Korea, 10-2024-0131645), which disclose 3D structure and property characterization techniques for nanocatalysts. The remaining authors declare no competing interests.

Additional information

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Supplementary Information

for

Atomic-scale 3D structural dynamics and functional degradation of Pt alloy nanocatalysts

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Supplementary Figure 1 | Experimental tomographic tilt series of the pristine PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for PtNi-pristine-p1 (**a**) and PtNipristine-p2 (**b**). Scale bar, 2 nm.

a	-71.0°	-67.4°	-63.4°	$-59.7°$	-56.3°	-53.1°	-50.2°	-47.5°	-45.0°
	-42.5°	-39.8°	-36.9°	-33.7°	-30.3°	-26.6°	-22.6°	-18.4°	-14.0°
	-9.5°	-4.8°	0.0°	4.8°	9.5°	14.0°	18.4°	22.6°	26.6°
	30.3°	33.7°	36.9°	39.8°	42.5°	45.0°	47.5°	50.2°	53.1°
	56.3°	59.7°	63.4°	67.4°	71.0°				
	-71.0°	-67.6°	-63.4°	-59.7°	-56.3°	-53.1°	-50.6°	-47.9°	-45.0°
$\mathbf{b}_{\mathbf{I}}$									
		-39.8°	-36.9°	-33.7°	-30.3°	-26.6°	-22.6°	-18.4°	-14.0°
	-42.3°								
	-9.5°	-5.6°	0.0°	4.4°	9.5°	14.0°	18.4°	22.6°	26.6°
	30.9°	33.9°	36.9°	39.6°	43.1°	45.0°	47.7°	51.0°	52.7°
	56.3°	59.7°	65.2°	67.4°	70.8°	74.1°			

Supplementary Figure 2 | Experimental tomographic tilt series of the 12k PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for PtNi-12k-p1 (**a**) and PtNi-12k-p2 (**b**). Scale bar, 2 nm.

Supplementary Figure 3 | Experimental tomographic tilt series of the 12k PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for PtNi-12k-p3 (**a**) and PtNi-12k-p4 (**b**). Scale bar, 2 nm.

Supplementary Figure 4 | Experimental tomographic tilt series of the 12k PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for PtNi-12k-p5 (**a**) and PtNi-12k-p6 (**b**). Scale bar, 2 nm.

Supplementary Figure 5 | Experimental tomographic tilt series of the pristine Ga-PtNi nanoparticles. a-c, Post-processed ADF-STEM images at different tilt angles (white text in each figure), for Ga-PtNi-pristine-p1 (**a**), Ga-PtNi-pristine-p2 (**b**), and Ga-PtNi-pristine-p3 (**c**). Scale bar, 2 nm.

Supplementary Figure 6 | Experimental tomographic tilt series of the 4k Ga-PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for Ga-PtNi-4k-p1 (**a**) and Ga-PtNi-4k-p2 (**b**). Scale bar, 2 nm.

Supplementary Figure 7 | Experimental tomographic tilt series of the 8k Ga-PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for Ga-PtNi-8k-p1 (**a**) and Ga-PtNi-8k-p2 (**b**). Scale bar, 2 nm.

Supplementary Figure 8 | Experimental tomographic tilt series of the 12k Ga-PtNi nanoparticles. a, **b**, Postprocessed ADF-STEM images at different tilt angles (white text in each figure), for Ga-PtNi-12k-p1 (**a**) and Ga-PtNi-12k-p2 (**b**). Scale bar, 2 nm.

Supplementary Figure 9 | Comparison of the zero-degree projections of PtNi nanoparticles during the AET experiments. a-h, The ADF-STEM images at zero-degrees, obtained at the beginning, middle, and end of the experiment, for PtNi-pristine-p1 (**a**), PtNi-pristine-p2 (**b**), PtNi-12k-p1 (**c**), PtNi-12k-p2 (**d**), PtNi-12k-p3 (**e**), PtNi-12k-p4 (**f**), PtNi-12k-p5 (**g**), and PtNi-12k-p6 (**h**). Scale bar, 2nm.

Supplementary Figure 10 | Comparison of the zero-degree projections of Ga-PtNi nanoparticles during the AET experiments. a-i, The ADF-STEM images at zero-degrees, obtained at the beginning, middle, and end of the experiment, for Ga-PtNi-pristine-p1 (**a**), Ga-PtNi-pristine-p2 (**b**), Ga-PtNi-pristine-p3 (**c**), Ga-PtNi-4k-p1 (**d**), Ga-PtNi-4k-p2 (**e**), Ga-PtNi-8k-p1 (**f**), Ga-PtNi-8k-p2 (**g**), Ga-PtNi-12k-p1 (**h**), and Ga-PtNi-12k-p2 (**i**). Scale bar, 2nm.

Supplementary Figure 11 | Experimentally determined 3D atomic structures of PtNi nanoparticles during potential cycles. a, Overall 3D atomic structure of PtNi (pristine and 12k) after given number of potential cycles. Note that these are the measured PtNi nanoparticles (PtNi-pristine-p1, PtNi-12k-p2, PtNi-12k-p3, PtNi-12k-p4, PtNi-12k-p5, PtNi-12k-p6) not illustrated in Fig. 1a-c. **b**, 3D atomic structures after a 180-degree rotation along the [001] axis from those in (**a**), with one octant of the nanoparticles removed to reveal the internal atomic structure. **c**, 1 Å thick internal slices at the center of the 3D tomograms perpendicular to the [001] direction. The intensity is shown in grayscale, with the atomic coordinates of Pt and Ni marked by blue and red dots, respectively. Scale bar, 2 nm.

Supplementary Figure 12 | Experimentally determined 3D atomic structures of Ga-PtNi nanoparticles during potential cycles. a, Overall 3D atomic structure of Ga-PtNi (pristine, 4k, 8k, and 12k) after given number of potential cycles. Note that these are the remaining Ga-PtNi nanoparticles (Ga-PtNi-pristine-p2, Ga-PtNipristine-p3, Ga-PtNi-4k-p2, Ga-PtNi-8k-p2, Ga-PtNi-12k-p1), excluding the representative nanoparticles specified in Fig. 1a-c. **b**, 3D atomic structures after a 180-degree rotation along the [001] axis from those in (**a**), with one octant of the nanoparticles removed to reveal the internal atomic structure. **c**, 1 Å thick internal slices at the center of the 3D tomograms perpendicular to the [001] direction. The intensity is shown in grayscale, with the atomic coordinates of Pt and Ni marked by blue and red dots, respectively. Scale bar, 2 nm.

Supplementary Figure 13 | Example experimental 2D projections where 3D octahedral-shaped nanoparticles appear spherical. a, **b**, Experimental projection images of Ga-PtNi-pristine-p1 at 60.2° (**a**) and Ga-PtNi-pristine-p2 at −65.5° (**b**), where the distinct octahedral shape evident in 3D tomogram is not clearly visible and appears spherical.

Supplementary Figure 14 | Architecture of the deep learning augmentation neural network. The neural network framework is based on a 3D U-Net¹. The boxes represent 3D feature maps. The black numbers below each feature map indicate the number of channels, and the pink numbers next to them denote the volume size of the 3D feature maps.

Supplementary Figure 15 | Learning curves during the training of the deep learning neural network. The blue dots and lines represent the root mean square error losses for the training dataset and red dots and lines represent the losses for the validation dataset.

^a These nanoparticles correspond to those shown in Fig. 1, while the remaining nanoparticles are presented in Supplementary Fig. 11.

	Ga-PtNi pristine			Ga-PtNi 4k		Ga-PtNi 8k		Ga-PtNi 12k	
	Ga-PtNi- pristine-p1 ^a	Ga-PtNi- pristine-p2	Ga-PtNi- pristine-p3	Ga-PtNi- $4k-p1a$	Ga-PtNi- 4k-p2	Ga-PtNi- 8k-p1	Ga-PtNi- 8k-p2 ^a	Ga-PtNi- 12k-p1	Ga-PtNi- $12k-p2^a$
STEM data acquisition									
Electron microscope type	TEAM 0.5 microscope at NCEM	FEI Titan Themis Z at INST	FEI Titan Themis Z at INST	TEAM 0.5 microscope at NCEM					
Acceleration voltage (kV)	200	300	300	200	200	200	200	200	200
Convergence semi-angle (mrad)	25.0	25.1	25.1	25.0	25.0	25.0	25.0	25.0	25.0
Detector inner angle (mrad)	43	41	41	43	43	43	43	43	43
Detector outer angle (mrad)	216	200	200	216	216	216	216	216	216
Pixel size (Å)	0.329	0.357	0.357	0.329	0.329	0.329	0.329	0.329	0.329
Screen current (pA)	10	10	10	10	10	10	10	10	10
Dwell time (µs)	3	3	3	3	3	3	3	3	3
# of consecutive images	4	4	4	4	4	4	4	$\overline{4}$	$\overline{4}$
# of projections	49	38	39	48	48	48	48	51	51
	-65.2	-73.8	-73.0	-66.3	-66.9	-64.5	-65.1	-69.9	-69.9
Tilt angle range (°)	$+63.7$	$+69.2$	$+73.0$	$+56.7$	$+57.1$	$+57.7$	$+57.0$	$+62.6$	$+63.1$
Electron dose (10 ⁵ e/ \AA ²)	3.4	2.2	2.3	3.3	3.3	3.3	3.3	3.5	3.5
3D reconstruction									
Algorithm	GENFIRE	GENFIRE	GENFIRE	GENFIRE	GENFIRE	GENFIRE	GENFIRE	GENFIRE	GENFIRE
Interpolation method	DFT	DFT	DFT	DFT	DFT	DFT	DFT	DFT	DFT
# of iterations	1000	1000	1000	1000	1000	1000	1000	1000	1000
Interpolation radius (pixel)	0.1	0.1	0.1	0.1	0.1	0.1	0.1	0.1	0.1
Oversampling ratio	$\overline{4}$	4	4	4	4	4	4	$\overline{4}$	4
Tracing result									
# of atoms									
Ni	11433	6810	7962	11578	7309	4267	10539	4116	4717
Pt	14714	9705	12022	14587	10142	7735	14451	7585	8768
B factor (A^2) Ni Pt	7.3 9.2	8.1 9.1	7.8 9.2	6.6 8.3	7.2 9.1	7.6 8.6	6.6 8.3	7.4 7.9	7.3 6.8
R factor $(\%)$	9.8	9.6	10.9	9.2	11.2	9.8	9.3	11.3	11.4
Precision estimation									
Standard deviation of Gaussian kernel (Å)	0.57	0.63	0.62	0.55	0.62	0.60	0.61	0.61	0.51
Accuracy of atom identification (%)	91.7	90.1	87.9	94.5	93.3	93.2	94.3	93.9	92.7
RMSD (pm)	38.6	39.2	41.1	34.8	33.1	30.4	34.2	32.1	31.5

Supplementary Table 2 | Experimental conditions, analysis parameters and tracing results of Ga-PtNi nanoparticles

^a These nanoparticles correspond to those shown in Fig. 1, while the remaining nanoparticles are presented in Supplementary Fig. 12.

Supplementary References

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