A note on an inversion algorithm for vertical ionograms for the prediction of plasma frequency profiles

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Building upon the concept of utilizing quasi-parabolic approximations to determine plasma frequency profiles from ionograms, we present a refined multi-quasi-parabolic method for modeling the E and F layers. While a recent study [AIP Advances 14, 065034 (2024)] [\[1\]](#page-4-0) introduced an approach in this direction, we identified several inaccuracies in its mathematical treatment and numerical results. By addressing these issues, we offer a clearer exposition and a more robust algorithm. Our method assumes a parabolic profile for the E layer and approximates the F layer with a series of concatenated quasi-parabolic segments, ensuring continuity and smoothness by matching derivatives at the junctions. Applied to daylight ionograms from the Jicamarca Observatory in Lima, our inversion algorithm demonstrates excellent agreement between the synthetic ionograms generated from our predicted plasma frequency profiles and the original measured data.

I. INTRODUCTION

Most of our knowledge about the ionosphere comes from ionogram records. These $h'(f)$ records give the apparent or virtual heights of reflection $h'(f)$ of a vertically transmitted radio wave, as a function of the wave frequency f . This paper aims to retrieve the electronic density profile from measured ionograms.

The analysis of ionograms consists basically of converting an observed $h'(f)$ curve, which gives the virtual height of reflection h' as a function of the wave frequency f, into an $N(h)$ curve giving the variation of the electron density N with height h . These two curves are related by

$$
h'(f) = \int_0^{h_r} \mu' dh,\tag{1}
$$

where the group refractive index μ' is a complicated function of f, N , and the strength and direction of the magnetic field. The height of reflection, h_r for the frequency f depends on f, N, and (for the extraordinary rays only) the strength of the magnetic field.

Previous efforts to solve this ill-posed problem have included lamination techniques, least-squares polynomial approximations, and ray tracing methods. Since there is no analytic solution to Eq. [\(1\)](#page-0-0)—that is, no direct expression for $N(h)$ in terms of $h'(f)$ —researchers have explored various numerical and approximation strategies. The lamination technique proposed by Reilly [\[4\]](#page-4-1) involves assuming various $N(h)$ model curves, passing them through a forward model to generate corresponding ionograms, and then comparing the resulting $h'(f)$ curves with those observed experimentally. In contrast, Reinisch et al. [\[3,](#page-4-2) [5\]](#page-4-3) utilized Chebyshev polynomial methods to approximate the F layer, aiming for a more efficient process in ionogram analysis. However, their software is proprietary, making it difficult to replicate and improve upon their work. More recently, Ankita et al. [\[6\]](#page-4-4) introduced a different approach by using electromagnetic wave propagation simulations based on Hamiltonian formulations for ray tracing.

In this study, we build upon the concepts introduced by Niu et al. [\[1\]](#page-4-0), who proposed using multivariate quasiparabolic layers to develop an inversion algorithm for approximating the plasma frequency profiles of the E and F ionospheric layers. While their underlying idea holds potential, we have identified several mathematical errors, mislabeled equations, and inconsistencies in their manuscript that make the methodology challenging to follow and replicate. To address these issues, we offer a clearer and more accurate exposition of the intended approach, providing detailed computations for transparency and understanding. This work is intended to be pedagogical, aiming to enhance comprehension of the inversion process. Additionally, recognizing the scarcity of open-source software for ionogram inversion, we are releasing our code to the community to facilitate further research and application.

An introduction for the QP layers and a treatment of how to use them to reconstruct the electron density profile is presented in Sec. [II.](#page-0-1) In Sec. [III](#page-2-0) we present a detailed inversion algorithm using ideas from the previous section. In Sec. [IV](#page-2-1) we describe the forward model, a series of detailed calculations to obtain the virtual heights given a plasma frequency profile. Finally, in Sec. [V](#page-3-0) we show some results of the predicted profiles based on a given ionograms.

II. MODEL: QP LAYERS

The quasi-parabolic model of the ionosphere is useful for representing electron density variations in layers such as the E and F regions; it was introduced by Forsterling et al. [\[8\]](#page-4-5). The electron density N_e as a function of the radial distance r from the Earth's center is described by:

$$
N_e = \begin{cases} N_m \left[1 - \left(\frac{r - r_m}{y_m}\right)^2 \right], & \text{for } r_b < r < r_m + y_m \\ 0, & \text{otherwise} \end{cases}
$$

where:

- N_e : Electron density, which attains a maximum value N_m at the center of the ionospheric layer.
- r: Radial distance from the Earth's center.
- r_m : Radial distance where the electron density reaches its peak value N_m .
- r_b : Radial distance at the base of the ionospheric layer, defined as $r_m - y_m$.
- y_m : Half-thickness of the ionospheric layer, controlling the layer's vertical extent.
- N_m : Maximum electron density at $r = r_m$.

This model confines the electron density within a parabolic profile, smoothly decaying to zero outside the bounds $r_b < r < r_m + y_m$. A very slight modification to the parabolic model permits the derivation of exact equations for ray-path parameters. This modified parabolic ionosphere will be termed the "quasi-parabolic" or, more simply, the "QP" model.

We make a QP layer the basic unit that will be used to model split sections of the F layer and the entirety of the E layer. We define a QP layer as

$$
f_{ni}^2 = f_{ci}^2 \left[1 \pm \left(\frac{r - r_{mi}}{y_{mi}} \right)^2 \left(\frac{b_i}{r} \right)^2 \right] \tag{2}
$$

where f_{ni} is the plasma frequency dependent on the real height r and f_{ci} is the critical frequency. However these physical meanings are lost once we consider the QP layer as a unit, and instead it is more a geometric construction. It can also be written as

$$
y_i = f_{ni}^2 = a_i \pm b_i \left(1 - \frac{r_{mi}}{r} \right)^2 \tag{3}
$$

where $y_i = f_{ni}^2$, $a_i = f_{ci}^2$, and $b_i = a_i \left(\frac{b_i}{y_{mi}}\right)^2$. There will be two kinds of QP layers: QP_i^+ and QP_i^- . Therefore they will be represented by equations y_i^+ and y_i^- , correspondingly. Each QP layer is parameterized by three numbers: a_i , b_i , and r_{mi} . See Fig [1.](#page-1-0)

The way we use Eq. [\(2\)](#page-1-1) to find the real heights, is solving the equation for a given frequency. That is, per QP layer, we select a bunch of plasma frequency points and solve the equation getting their corresponding real heights.

A. Data

We assume we have the autoscaled ionogram virtual heights for frequencies $f_i \in [0, f_F]$ with arbitrary strides between frequency points, and also the frequency position of the E layer critical frequency f_E . Big gaps between frequency data points may cause undesired behaviour in the algorithm.

B. E layer

It is well known that the E-layer virtual heights per frequency f can be modeled as

$$
h'(f) = r_b + \frac{1}{2} y_m \frac{f}{f_E} \ln\left(\frac{f_E + f}{f_E - f}\right) \tag{4}
$$

Where r_b is the height where the ionogram starts. In the case of E layers, we usually don't have the complete trace so it is necessary to search for r_b and y_m values. If we were to use a QP model for the E layer instead of Eq. [\(4\)](#page-1-2), we would still need to compute the same values since we take only f_E , the E layer critical frequency, for granted. We use a brute-force approach to find the best parameters for the E-layer using a simple two for-loop. For each pair (r_{bE}, y_{mE}) we compare the produced virtual heights $h'(f)$ using the Eq. [\(4\)](#page-1-2), or with the forward model in case we use a QP model, versus the original ionogram. Finally, the probe pair with the least error difference is the one we take. The metric used for measuring the error we use is the root-mean-squared error.

Because the E layer presents a pronounced steep, we consider a probe frequency points that are not evenly distributed in, say, a range $[A, B]$, but more densily distributed near the upper bound B . We use the following equation to obtain a suitable distribution of values

$$
x_i = A + (B - A) \frac{1 - \exp(-k\frac{i}{N})}{1 - \exp(-k)}.
$$
 (5)

C. QP layers concatenation for F layer reconstruction

Next, we approximate the F layer by concatenating several QP layers by alternating between QP^- to QP^+

FIG. 1. (Color online) Description of what the concatenation of QP layers aims to build in a plasma frequency vs real height plane. The dotted black line would be the plasma frequency profile.

and QP^+ to QP^- layers. We start the F layer with a QP^+ layer. We ensure the continuity of the plasma frequency profile between QP_i and QP_{i-1} by making its derivatives equal at r_i , the real height corresponding to last frequency point we considered in the QP_{i-1} layer. See Fig. [1.](#page-1-0)

$$
\left. \frac{dy_i^{\pm}}{dr} \right|_{r=r_i} = \left. \frac{dy_{i-1}^{\mp}}{dr} \right|_{r=r_i} \quad \text{or} \quad \left. \frac{dy_i^{\mp}}{dr} \right|_{r=r_i} = \left. \frac{dy_{i-1}^{\pm}}{dr} \right|_{r=r_i},\tag{6}
$$

From these relations we can get expressions for b_i and r_{mi} , dependent on a_i and the parameters of the QP_{i-1} layer. See Eq. [\(3\)](#page-1-3)

Let us compute the derivatives for the different types of QP layers. For y_i^-

$$
y_i^-(r) = a_i - b_i \left(1 - \frac{r_{mi}}{r}\right)^2
$$
 (7)

its corresponding derivative is

$$
y_i^-(r)' = -\frac{2b_i r_{mi}}{r^2} \left(1 - \frac{r_{mi}}{r}\right)
$$
 (8)

On the other hand, for y^+

$$
y_i^+(r) = a_i + b_i \left(1 - \frac{r_{mi}}{r}\right)^2 \tag{9}
$$

its derivative is

$$
y_i^+(r)' = \frac{2b_i r_{mi}}{r^2} \left(1 - \frac{r_{mi}}{r}\right) \tag{10}
$$

Then, as we mentioned, there will be two cases: QP^- to QP^+ or QP^+ to QP^- . First, for the case of QP_{i-1}^- to QP_i^+ , to find the dependence of parameters b_i and r_{mi} with the ones of QP_{i-1} and ensure the continuity of the curve, we equate $y_{i-1}^-(r_i)' = y_i^+(r_i)'$. r_i is the point where both curves meet, computationally speaking r_i is the last point we took from the curve QP_{i-1}^- . We find

$$
r_{mi} = \frac{r_i^2 y_{i-1}^-(r_i)'}{2(y_{i-1}(r_i) - a_i) + r_i y_{i-1}^-(r_i)'}\tag{11}
$$

$$
b_i = \frac{\left[2(y_{i-1}(r_i) - a_i) + r_i y_{i-1}^-(r_i)'\right]^2}{4(y_{i-1}(r_i) - a_i)}\tag{12}
$$

Finally, for the case of QP_{i-1}^+ to QP_i^- , we equate

$$
y_{i-1}^+(r_i)' = y_i^-(r_i)'.\tag{13}
$$

$$
r_{mi} = \frac{r_i^2 y_{i-1}^+(r_i)'}{2(y_{i-1}^+(r_i) - a_i) + r_i y_{i-1}^+(r_i)'}\tag{14}
$$

$$
b_i = -\frac{\left[2(y_{i-1}^+(r_i) - a_i) + r_i y_{i-1}^+(r_i)'\right]^2}{4(y_{i-1}^+(r_i) - a_i)}\tag{15}
$$

III. ALGORITHM

The general idea is that for each QP layer we will try different values of f_{ci} , each giving a its own set of parameters, which at the same time depend on values computed in the previous QP_{i-1} . Next, for each f_{ci} , we use a batch of N plasma frequency points for which we calculate their real heights and subsequently their virtual heights using the forward model. We compare them against the measured virtual heights and keep the batch with the least error. Finally, we append the best one and move on onto the next layer.

A more detailed explanation is as follows: the algorithm starts with the calculation of the E-layer, also known as QP_0^- layer. Then, we can iteratively compute the next QP_i parameters since we already know the ones from the QP_{i-1} layer meaning r_{mi} and b_i will be left as a function of $a_i = f_{ci}^2$. To find f_{ci} , we use an exhaustive search algorithm, which means trying several values of f_{ci} by brute force. If the QP_i layer is a QP^+ layer, we will search f_{ci} in the range $[f_L + \epsilon, f_L + 2.0]$, and, if, on the other hand, we are computing a QP^- layer, the range will be $[f_L - 2.0, f_L - \epsilon]$. For each of these f_{ci} values, we take a batch of N plasma frequency points including f_{ci} and compute their real and virtual heights with the help of the forward mode and compare them against the real measurements from the ionogram. The batch producing the least error is attached to our fp profile array. Once we are done with the QP_i layer, we consider the last element of the fp profile to be the intersection point between the QP_i and QP_{i+1} layers. Finally, we repeat the process. The number of QP layers should be defined beforehand.

IV. FORWARD MODEL

Each time we add new plasma frequency and real height data points to our plasma frequency profile curve, we need to test it against the original ionogram and calculate the error. To achieve this we need a forward model that can calculate an ionogram given a plasma frequency profile.

Another way to express the ionospheric virtual height reflection is

$$
h(f) = \int_0^{z_r} \frac{dz}{n(z)}\tag{16}
$$

FIG. 2. Close up to the $n(z)^2$ function.

and by looking at the expression of $n(z) = \sqrt{1 - \frac{f_p(z)^2}{f_z^2}}$ $\frac{f(z)^2}{f^2},$ plus the condition that $f_p(z_r) = f$, we can foresee that the curve of $n(z)^2$ will look like an inverse function beginning in the coordinate $(0, 1)$ and then converging towards $n(z)^2 = 0$ at $z = z_r$. If we zoom in on a curve segment, the ending points would be like the ones in Fig. [2.](#page-3-1) Using triangle relations we get

$$
\frac{z_{i+1} - z_i}{n^2(z_i) - n^2(z_{i+1})} = \frac{z - z_i}{n^2(z_i) - n^2(z)}\tag{17}
$$

where $n(z)$ can be extracted from this equation

$$
n^{2}(z) = n^{2}(z_{i}) - (z - z_{i}) \left[\frac{n^{2}(z_{i}) - n^{2}(z_{i+1})}{z_{i+1} - z_{i}} \right]
$$
 (18)

and replace it in Eq. [16.](#page-2-2) Because the integrand is an inverse function the biggest contributions will come from the segment with z_i value closer to z_r . If we treat the integral from Eq. [16](#page-2-2) by parts (imagine our line of integration is formed by N points), then there will be two different contributions, one coming from the last segment (z_1, z_r) and another from (z_0, z_1) .

$$
h(f) = \int_{z_1}^{z_r} \frac{dz}{n(z)} + \int_0^{z_1} \frac{dz}{n(z)}
$$
(19)

First, we deal with the latter.

$$
= \sum_{i=0}^{N-2} \int_{z_i}^{z_{i+1}} \frac{dz}{n(z)}
$$

=
$$
\sum_{i=0}^{N-2} \int_{z_i}^{z_{i+1}} \frac{1}{\sqrt{n^2(z_i) - (z-z_i) \left[\frac{n^2(z_i) - n^2(z_{i+1})}{z_{i+1} - z_i} \right]}} dz
$$
(20)

$$
= \int_{0}^{z_{1}} \frac{dz}{n(z)}
$$

=
$$
\sum_{i=0}^{N-2} \int_{0}^{z_{i+1}-z_{i}} \frac{1}{\sqrt{n^{2}(z_{i}) - z' \left[\frac{n^{2}(z_{i}) - n^{2}(z_{i+1})}{z_{i+1} - z_{i}}\right]}} dz',
$$
 (21)

Then using the result $\int \frac{1}{\sqrt{a-bx}} dx = -\frac{2\sqrt{a-bx}}{b} + C$ to calculate the integral analytically

$$
\int_0^{z_1} \frac{dz}{n(z)} = \sum_{i=0}^{N-2} \frac{2(z_{i+1} - z_i)}{\sqrt{n^2(z_i)} + \sqrt{n^2(z_{i+1})}}
$$

=
$$
\sum_{i=0}^{N-2} \frac{2\Delta z}{n(z_i) + n(z_{i+1})}
$$
(22)

Secondly, we deal with the first integral in Eq. [\(19\)](#page-3-2). As the upper limit for the integral is z_r , we need to remember that $n(z = z_r) = 0$, then because of $z_r =$ $\frac{n^2(z_i)}{n^2(z_i)-n^2(z_{i+1})}(z_{i+1}-z_i)+z_i$ the integral will be

$$
n^{2}(z_{i}) \int_{0}^{\frac{z_{i+1}-z_{i}}{n^{2}(z_{i})-n^{2}(z_{i+1})}} \frac{dz}{n(z')} \tag{23}
$$

already applied the change of variables. The result is

$$
\int_{z_i}^{z_r} \frac{dz}{n(z)} = 2n(z_i) \frac{(z_{i+1} - z_i)}{n^2(z_i) - n^2(z_{i+1})}
$$
(24)

with $z_i = z_{N-1}$ and $z_{i+1} = z_N$. Then, by calculating both analytically both integrals we can get an accurate value for the virtual height given a plasma frequency profile.

V. RESULTS

To test our inverse algorithm and forward model, we used data measured in the Jicamarca Observatory in Lima, Peru. We concentrate in daylight ionograms as they present some data points for the E layer. We do not support ionograms with only the F layer present.

The original measured ionograms will be represented by magenta dots, the reconstructed plasma frequency profile in a yellow solid line, and the synthetic ionogram calculated from the reconstructed plasma frequency profile will be represented by empty black squares. For the sake of brevity, we present results with an hour of difference during daylight.

Even though the final reconstructed ionogram does not perfectly fit the original one, see Fig. [9](#page-7-0) , in general shows good agreement, see Figs. [4](#page-5-0) [5](#page-5-1) [6](#page-6-0) [8.](#page-6-1) But more importantly, the results give us a good idea of what the plasma frequency or electron density profile look like.

To obtain this results, we usually computed between 25 to 30 QP layers with 4 data points per layer. However, to stabilize the virtual heights of the first produced synthetic ionogram data points in the F layer, we used a single data point for the first 10 QP layers. Additionally, our algorithm is able to complete the missing data points in the E layer, see for example Fig. [9.](#page-7-0)

The open-source code and data used for this research

- [1] L. Niu, L. Wen, C. Zhou, and M. Deng, "A profile inversion method for vertical ionograms," AIP Advances, vol. 14, no. 6, p. 065034, Jun. 2024. doi: 10.1063/5.0208687.
- [2] J. E. Titheridge, "A new method for the analysis of ionospheric h'(f) records," Journal of Atmospheric and Terrestrial Physics, vol. 21, no. 1, pp. 1-12, 1961. doi: 10.1016/0021-9169(61)90185-4.
- [3] B. W. Reinisch and X. Huang, "Automatic calculation of electron density profiles from digital ionograms: 3. Processing of bottomside ionograms," Radio Science, vol. 18, no. 3, pp. 477-492, 1983. doi: 10.1029/RS018i003p00477.
- [4] M. H. Reilly and J. D. Kolesar, "A method for real height analysis of oblique ionograms," Radio Science, vol. 24, no. 04, pp. 575-583, 1989. doi: 10.1029/RS024i004p00575.
- [5] J. E. Titheridge, "Direct Manual Calculations of Ionospheric Parameters Using a Single-Polynomial Analysis," Radio Science, vol. 2, no. 10, pp. 1237-1253, 1967. doi:

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- [6] M. Ankita, S. Tulasi Ram, "Iterative Gradient Correction (IGC) Method for True Height Analysis of Ionograms," Radio Science, vol. 58, Nov. 2023. doi: 10.1029/2023RS007808.
- [7] K. G. Budden, "The Numerical Solution of the Differential Equations Governing the Reflexion of Long Radio Waves from the Ionosphere. II," Philosophical Transactions of the Royal Society of London. Series A, Mathematical and Physical Sciences, vol. 248, no. 939, pp. 45-72, 1955. [Online]. Available: [http://www.jstor.org/stable/91622.](http://www.jstor.org/stable/91622)
- [8] V. K. Forsterling and H. Lassen, "Die Ionisation der Atmosphäre und die Ausbreitung der kurzen elektrischen Wellen (IQ-100 m) über die Erde. III," Zeitschrift für technische Physik, vol. 12, pp. 502-527, 1931.

FIG. 3.

FIG. 4.

FIG. 5.

FIG. 6.

FIG. 7.

FIG. 8.

FIG. 9.