

ATIYAH-PATODI-SINGER BOUNDARY CONDITION AND A SPLITTING FORMULA OF A SPECTRAL FLOW

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ABSTRACT. We describe a relation between Atiyah-Patodi-Singer boundary condition and a global elliptic boundary condition which naturally appears in formulating a splitting formula for a spectral flow, when we decompose the manifold into two components. Then we give a variant of the splitting formula with the Hörmander index as a correction term.

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1. INTRODUCTION

In the paper [FO] we formulated and proved a splitting formula of a spectral flow for a continuous family of first order selfadjoint elliptic differential operators $\{A_t\}_{t \in [0,1]}$ defined on a closed manifold M . This is an addition formula of a spectral flow when we decompose a manifold into two components along a hypersurface Σ , $M = M_- \cup_{\Sigma} M_+$, $\partial M_{\pm} = \Sigma$. For such a family $\{A_t\}$ considered on the whole closed manifold M , an integer, called spectral flow, is well defined, and to “observe” this quantity we cut the manifold along a hypersurface, then we can “observe” from the hypersurface a quantity “Maslov index” which is a curve of boundary data of solutions of operators. This quantity can be understood as the spectral flow. This is just a spectral flow formula ([Yo], [Ni]) where manifolds need not be separated into two components. If the manifold is separated into two parts by the hypersurface, then we will have two Maslov indexes which together give the whole information of the spectral flow. For this observation we must make clear which family we are observing, i.e., to get a family of selfadjoint Fredholm extensions we must impose a suitable elliptic boundary condition on the family $\{A_t\}$ when we restrict the operators on each component M_{\pm} . This condition appears in a natural way in our formulation to write down the splitting formula and reflects the influence from one side to other side. On the other hand Atiyah-Patodi-Singer boundary condition is described

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based on the boundary data. Nevertheless these two are relating each other in the case of the operators of the product form near the hypersurface.

The purpose of this paper is to describe a relation between our global elliptic boundary condition and Atiyah-Patodi-Singer boundary condition. There are several such splitting formulas ([CLM2], [DK], [KL],[Ta],[Ni]) and the boundary condition treated there is mostly Atiyah-Patodi-Singer boundary condition. So the result in this paper will give us a sight of our global elliptic boundary condition for operators satisfying suitable analytic assumptions.

We follow the theory of symplectic Hilbert spaces, especially of the Fredholm-Lagrangian-Grassmannian and the Maslov index in the infinite dimension which were discussed in the papers [Fu] and [FO] precisely(also see [Go], [CLM1] and [RS]).

In §2 we explain a global elliptic boundary condition appearing in the splitting formula for a spectral flow and state a relation between it and Atiyah-Patodi-Singer boundary condition in terms of the Fredholm-Lagrangian-Grassmannian.

In §3, first we recall a L_2 -reduction theorem [Fu] and by applying this we give a proof of Theorem (2.7).

In §4 as an application of Theorem (2.7) we give a variant of a splitting formula of a spectral flow with Hörmander index as a correction term.

2. A GLOBAL ELLIPTIC BOUNDARY CONDITION

In this section we explain an elliptic boundary condition we introduced in [FO].

Let A be a first order selfadjoint elliptic differential operator defined on a real vector bundle \mathbb{E} on a closed manifold M . Let Σ be a hypersurface of M along which M is separated into two components M_{\pm} , $M = M_- \cup_{\Sigma} M_+$, $\partial M_{\pm} = \Sigma$, and we denote the first order Sobolev space on M (resp. M_{\pm}) taking values in the real vector bundle \mathbb{E} by $H^1(M, \mathbb{E})$ (resp. $H^1(M_{\pm}, \mathbb{E}|_{M_{\pm}})$). For the subspace in $H^1(M_{\pm}, \mathbb{E}|_{M_{\pm}})$ with vanishing boundary values we denote it by $H_0^1(M_{\pm}, \mathbb{E}|_{M_{\pm}})$. These are the domains of the minimal closed extensions of the operators A considered on $C_0^{\infty}(M_{\pm} \setminus \Sigma, \mathbb{E}|_{M_{\pm} \setminus \Sigma})$ and we denote them by $\mathfrak{D}_{\mathbf{m}}^{\pm} = H_0^1(M_{\pm}, \mathbb{E}|_{M_{\pm}})$. Then we denote by \mathcal{A}_{\pm}^* their adjoint operators considered on M_{\pm} and by $\mathfrak{D}_{\mathbf{M}}^{\pm}$ their domains of definitions, i.e., $f \in \mathfrak{D}_{\mathbf{M}}^{\pm}$ if $f \in L_2(M_{\pm}, \mathbb{E}|_{M_{\pm}})$ and $A(f) \in L_2(M_{\pm}, \mathbb{E}|_{M_{\pm}})$ in the distribution sense.

We must put two assumptions **(a1)** and **(a2)** on the selfadjoint elliptic operator A :

(a1): A satisfies the *unique continuation property* with respect to the hypersurface Σ , that is,

$$(2.1) \quad \text{Ker}(\mathcal{A}_{\pm}^*) \cap \mathfrak{D}_{\mathbf{m}}^{\pm} = \{0\}.$$

(a2): On a tubular neighborhood $\mathcal{N} \cong (-1, 1) \times \Sigma$ the operator A is of the product form, that is,

$$(2.2) \quad A = \sigma \left(\frac{\partial}{\partial u} + B_0 \right),$$

where σ is a bundle map on $\mathbb{E}|_{\mathcal{N}}$ which does not depend on the normal variable $u \in (-1, 1)$, the operator B_0 is a selfadjoint elliptic operator on Σ and also does not depend on the normal variable u .

We identify $\mathcal{N} \cap M_- \cong (-1, 0] \times \Sigma$ and $\mathcal{N} \cap M_+ \cong [0, 1) \times \Sigma$.

The selfadjointness of the operator A implies that

$$(2.3) \quad \sigma^2 = -\text{Id}, \quad {}^t\sigma = -\sigma$$

$$(2.4) \quad \sigma \circ B_0 + B_0 \circ \sigma = 0,$$

where the transpose ${}^t\sigma$ is taken with respect to a suitable metric on the vector bundle \mathbb{E} . We denote the inner product on the L_2 sections of \mathbb{E} by $\langle \cdot, \cdot \rangle$. Then σ defines an almost complex structure and a compatible symplectic structure on $L_2(\Sigma, \mathbb{E}|_\Sigma) = L_2(\Sigma)$.

Let \mathfrak{D}_0 be a subspace in $H^1(M_-, \mathbb{E}|_{M_-})$ defined by

$$\mathfrak{D}_0 = \{f \in H^1(M_-, \mathbb{E}|_{M_-}) \mid \exists \tilde{f} \in H^1(M, \mathbb{E}) \text{ such that } \tilde{f}|_{M_-} = f, \text{ and } A(f) = 0 \text{ on } M_+\}.$$

Hereafter we will state properties only for the case M_- , but shall use the corresponding results for M_+ , if necessary.

We denote the restriction of \mathcal{A}_-^* to \mathfrak{D}_0 by $A_{\mathfrak{D}_0}$, then under the two assumptions **(a1)** and **(a2)** we have

Proposition 2.1. *The operator $A_{\mathfrak{D}_0}$ satisfies the inequality:*

$$(2.5) \quad \|f\|_1 \leq c(\|A_{\mathfrak{D}_0}(f)\| + \|f\|), \quad \forall f \in \mathfrak{D}_0$$

with a positive constant $c > 0$, where $\|f\|_1$ denotes the first order Sobolev norm. And so $A_{\mathfrak{D}_0}$ is selfadjoint and has compact resolvents.

This property is basic to state Theorem (2.7) below. We already made use of this property in our paper ([FO]). The proof is given upon L_2 -reduction theorem of the Maslov index in the infinite dimension. In the next section we explain some part of a proof of this proposition, and together with the help of Rellich's Theorem we prove Theorem (2.7).

Remark 2.2. If A is invertible on $H^1(M, \mathbb{E})$, then the same holds for $A_{\mathfrak{D}_0}$ as in the above proposition without assumptions **(a1)** and **(a2)**, because we have

$$\|f\|_{H^1(M_-, \mathbb{E}|_{M_-})} \leq C(\|f\|_{H^1(M, \mathbb{E})} + \|f\|_{L_2(M)}) = C(\|f\|_{H^1(M_-, \mathbb{E}|_{M_-})} + \|f\|_{L_2(M)}) \leq C'\|f\|_{H^1(M_-, \mathbb{E}|_{M_-})}$$

for $f \in H^1(M, \mathbb{E})$ satisfying $A(f) = 0$ on M_+ .

Let $\{\ell_k\}_{k \in \mathbb{Z} \setminus \{0\}}$ be the eigenvalues of the operator B_0 and we denote corresponding orthonormal eigensections by $\{\varphi_k\}$. From the properties (2.3) and (2.4) we have $\ell_k = -\ell_{-k} > 0$ for $k = N_0 + 1, N_0 + 2, \dots$ with $N_0 = 1/2 \cdot \dim \text{Ker}(B_0)$ (see Remark (2.4) below) and $\ell_k = 0$ for $0 < |k| \leq N_0$.

For a section φ on Σ let

$$\varphi = \sum_{k \in \mathbb{Z} \setminus \{0\}} a_k \varphi_k$$

be the eigensection-expansion, then the Sobolev space $H^s(\Sigma, \mathbb{E}|_\Sigma)$ of order $s \in \mathbb{R}$ on Σ is characterized as

$$H^s(\Sigma, \mathbb{E}|_\Sigma) = \left\{ \varphi = \sum_{k \in \mathbb{Z} \setminus \{0\}} a_k \varphi_k \mid \sum_{k \in \mathbb{Z} \setminus \{0\}} |a_k|^2 |\ell_k|^{2s} < \infty \right\}.$$

Let \mathfrak{D}_{APS}^0 be a subspace in $H^1(M_-, \mathbb{E}|_{M_-})$ such that

$$(2.6) \quad \mathfrak{D}_{APS}^0 = \left\{ f \in H^1(M_-, \mathbb{E}|_{M_-}) \mid \text{if } f|_\Sigma = \sum_{k \in \mathbb{Z} \setminus \{0\}} a_k \varphi_k, \text{ then } a_{-k} = 0 \text{ for } k = 1, 2, 3, \dots \right\},$$

and we denote the restriction of \mathcal{A}_-^* to \mathfrak{D}_{APS}^0 by $A_{\mathfrak{D}_{APS}^0}$. This non-local boundary condition (2.6) is called ‘‘Atiyah-Patodi-Singer boundary condition’’. Then

Proposition 2.3. ([APS]) *The operator $A_{\mathfrak{D}_{APS}^0}$ defined on \mathfrak{D}_{APS}^0 is a selfadjoint operator with compact resolvents.*

Now let $\beta^- = \mathfrak{D}_{\mathbf{M}}^-/\mathfrak{D}_{\mathbf{m}}^-$ be the space of boundary values. Here the maximal domain $\mathfrak{D}_{\mathbf{M}}^-$ is equipped with the norm $\|\cdot\|^G$ given by the graph inner product:

$$\langle f, g \rangle^G = \langle f, g \rangle + \langle \mathcal{A}_-^*(f), \mathcal{A}_-^*(g) \rangle.$$

The space β^- has a structure of a symplectic Hilbert space with the symplectic form

$$(2.7) \quad \omega([f], [g]) = \langle \mathcal{A}_-^*(f), g \rangle - \langle f, \mathcal{A}_-^*(g) \rangle, \quad f, g \in \mathfrak{D}_{\mathbf{M}}^-,$$

and is realized in the distribution space on Σ :

$$\begin{aligned} \beta^- &= \left\{ f \in H^{-1/2}(\Sigma, \mathbb{E}_{|\Sigma}) \mid f = \sum_{k \in \mathbb{Z} \setminus \{0\}} c_k \varphi_k, \text{ with } \sum_{k > 0} |c_k|^2 \ell_k < \infty \text{ and} \right. \\ &\quad \left. \sum_{k < 0} |c_k|^2 |\ell_k|^{-1} < \infty \right\} \\ &= \left\{ f \in H^{1/2}(\Sigma, \mathbb{E}_{|\Sigma}) \mid f = \sum_{k > 0} c_k \varphi_k, \text{ with } \sum_{k > 0} |c_k|^2 \ell_k < \infty \right\} \\ &+ \left\{ f \in H^{-1/2}(\Sigma, \mathbb{E}_{|\Sigma}) \mid f = \sum_{k < 0} c_k \varphi_k, \text{ with } \sum_{k < 0} |c_k|^2 |\ell_k|^{-1} < \infty \right\} \\ &= \theta_+^- + \theta_-^-. \end{aligned}$$

For the determination of the space β^- see [Ho1] and [APS].

Remark 2.4. By the conditions **(a1)** and **(a2)** we know that $\text{Ker}(B_0)$ is a finite dimensional symplectic subspace of β^- , so that we choose eigensections $\{\varphi_k\}$ for $0 < |k| \leq N_0$ in such a way that the subspaces spanned by $\{\varphi_k\}_{-N_0 \leq k < 0}$ and $\{\varphi_k\}_{0 < k \leq N_0}$ are mutually transversal Lagrangian subspaces in $\text{Ker}(B_0)$.

Remark 2.5. Of course for smooth sections of $\mathbb{E}_{|\Sigma}$ (and also for L_2 -sections) the symplectic structure defined by σ coincides with ω defined in (2.7).

Let $\gamma^- : \mathfrak{D}_{\mathbf{M}}^- \rightarrow \beta^-$ be the projection map, then the image $\gamma^-(\text{Ker}(\mathcal{A}_-^*))$ is a Lagrangian subspace and the pairs $(\gamma^-(\text{Ker}(\mathcal{A}_-^*)), \gamma^-(\mathfrak{D}_{APS}^0))$ and $(\gamma^-(\text{Ker}(\mathcal{A}_-^*)), \gamma^-(\mathfrak{D}_0))$ are Fredholm pairs. The spaces $\gamma^-(\text{Ker}(\mathcal{A}_-^*))$ is called *Cauchy data space*. Further, for each Lagrangian subspace $\lambda \subset \beta^-$ the operator $\mathcal{A}_-^*|_{(\gamma^-)^{-1}(\lambda)}$ is a selfadjoint realization and if $(\lambda, \gamma^-(\text{Ker}(\mathcal{A}_-^*)))$ is a Fredholm pair, then $\mathcal{A}_-^*|_{(\gamma^-)^{-1}(\lambda)}$ is a selfadjoint Fredholm operator.

Remark 2.6. Note that the Lagrangian property of the Cauchy data space $\gamma^-(\text{Ker}(\mathcal{A}_-^*))$ would not be trivial. To prove this property we rely on the existence of at least one selfadjoint realization of $A|_{\mathfrak{D}_{\mathbf{m}}^-}$ (= restriction of \mathcal{A}_-^* to a suitable subspace in $\mathfrak{D}_{\mathbf{M}}^-$) with compact resolvents (or a unbounded selfadjoint Fredholm extension). For our case this realization is given by the operator $A_{\mathfrak{D}_{APS}^0}$, and for which proof we use the assumption **(a2)**.

We denote by $\mathcal{F}\Lambda_\lambda(\beta^-)$ the space of Lagrangian subspaces μ in β^- such that (μ, λ) is a Fredholm pair:

$$\mathcal{F}\Lambda_\lambda(\beta^-) = \{\mu \subset \beta^- \mid \mu \text{ is a Lagrangian subspace and } (\mu, \lambda) \text{ is a Fredholm pair}\}.$$

Now we have

Theorem 2.7.

$$(2.8) \quad \mathcal{F}\Lambda_{\gamma^-(\mathfrak{D}_0)}(\beta^-) = \mathcal{F}\Lambda_{\gamma^-(\mathfrak{D}_{APS}^0)}(\beta^-),$$

more precisely the orthogonal projection operators onto the subspace $\gamma^-(\mathfrak{D}_0)$ and that onto the subspace $\gamma^-(\mathfrak{D}_{APS}^0) = \theta_+^-$ differ by a compact operator.

We prove this in the next section.

Remark 2.8. Let H be a symplectic Hilbert space and we regard H as a complexification of a Lagrangian subspace λ . Then a Lagrangian subspace μ is of the form $U(\lambda) = \mu$ with a unitary operator U of the form $\text{Id} + \text{compact operator}$, then $\mathcal{F}\Lambda_\lambda(H) = \mathcal{F}\Lambda_\mu(H)$. Also this property is equivalent to the condition that the difference of the orthogonal projection operators onto the Lagrangian subspaces λ and μ is compact. For such two Lagrangian subspaces λ and μ and two arbitrary Lagrangian subspaces ν_0 and ν_1 in $\mathcal{F}\Lambda_\lambda(H) = \mathcal{F}\Lambda_\mu(H)$ we have a well-defined integer $\sigma(\nu_0, \nu_1; \lambda, \mu)$, called Hörmander index. This is the difference of the Maslov indexes

$$\sigma_H(\nu_0, \nu_1; \lambda, \mu) = \mathbf{Mas}(\{c(t)\}, \lambda) - \mathbf{Mas}(\{c(t)\}, \mu),$$

where the path $\{c(t)\}$ is in $\mathcal{F}\Lambda_\lambda(H)$ connecting ν_0 and ν_1 and the difference does not depend on any such paths. Here the Maslov index $\mathbf{Mas}(\{c_t\}, \lambda)$ is, in a sense, the intersection number with the ‘‘Maslov cycle’’ $\mathfrak{M}_\lambda = \{\mu \in \mathcal{F}\Lambda_\lambda(H) \mid \mu \cap \lambda \neq \{0\}\}$ ([Ho2], [Fu]).

3. SYMPLECTIC REDUCTION THEOREM

In this section after recalling a symplectic reduction theorem ([Fu]) we prove Theorem (2.7).

Let $(\mathcal{B}, \omega_{\mathcal{B}})$ and (L, ω_L) be two symplectic Hilbert spaces ($\omega_{\mathcal{B}}$ is the symplectic form and so on) with decompositions by Lagrangian subspaces θ_-, θ_+, L_- and L_+ (*Polarized symplectic Hilbert space*):

$$(3.1) \quad \mathcal{B} = \theta_- + \theta_+, \quad L = L_- + L_+.$$

We assume that there are continuous injective maps $\mathbf{i}_+ : L_+ \rightarrow \theta_+$ and $\mathbf{i}_- : \theta_- \rightarrow L_-$ having dense images such that

$$(3.2) \quad \omega_{\mathcal{B}}(\mathbf{i}_+(a), x) = \omega_L(a, \mathbf{i}_-(x)) \text{ for any } x \in \theta_- \text{ and } a \in L_+.$$

Then

Proposition 3.1. ([Fu]) *There is a continuous map $\tau : \mathcal{F}\Lambda_{\theta_-}(\mathcal{B}) \rightarrow \mathcal{F}\Lambda_{L_-}(L)$ such that for any continuous curve $\{c(t)\}_{t \in [0,1]}$ in $\mathcal{F}\Lambda_{\theta_-}(\mathcal{B})$*

$$\mathbf{Mas}(\{c(t)\}, \theta_-) = \mathbf{Mas}(\{\tau(c(t))\}, L_-).$$

The map τ is defined in the following way:

$$\tau(\nu) = \{b + a \in L = L_+ + L_- \mid \exists x \in \theta_- \text{ such that } \mathbf{i}_+(b) + x \in \nu \text{ and } a = \mathbf{i}_-(x)\}.$$

For any decomposition of $\theta_- = F + F'$ by closed subspaces F ($\dim F < +\infty$) and F' , we can decompose L_+ by closed subspaces in such a way that $L_+ = G + G'$ with $\dim G = \dim F$ and $G + \mathbf{i}_-(F)$ is a symplectic subspace in L . Also in this case the subspace $F + \mathbf{i}_+(G)$ is a symplectic subspace in \mathcal{B} . Moreover the subspaces $F + \overline{\mathbf{i}_+(G')}$ and $G + \mathbf{i}_-(F')$ are Lagrangian subspaces.

Then by replacing θ_- with $F' + \mathbf{i}_+(G)$, θ_+ with $F + \overline{\mathbf{i}_+(G')}$, L_- with $G + \overline{\mathbf{i}_+(F')}$ and L_+ with $\mathbf{i}_-(F) + G'$ and also by replacing the maps \mathbf{i}_\pm in an obvious way we have a similar situation as in (3.1) and (3.2). We shall denote these new maps by $\tilde{\mathbf{i}}_\pm$, although the resulting maps τ between Fredholm-Lagrangian-Grassmannians $\mathcal{F}\Lambda_{\theta_-}(\mathcal{B}) = \mathcal{F}\Lambda_{F'+\mathbf{i}_+(G)}(\mathcal{B})$ and $\mathcal{F}\Lambda_{L_-}(L) = \mathcal{F}\Lambda_{G+\overline{\mathbf{i}_+(F')}}(L)$ coincides.

Note that the arguments above are guaranteed that the spaces \mathcal{B} and L are Hilbert spaces (see [KS] for symplectic Banach spaces).

We apply this proposition to the case $\mathcal{B} = \beta^+ = \theta_+^+ + \theta_+^-$ and $L = \beta^- = \theta_+^- + \theta_-^-$. Note that the space β^+ is defined as follows:

$$\begin{aligned} \beta^+ &= \{f \in H^{-1/2}(\Sigma, \mathbb{E}_\Sigma) \mid f = \sum_{k \in \mathbb{Z} \setminus \{0\}} c_k \varphi_k, \text{ with } \sum_{k>0} |c_k|^2 \ell_k^{-1} < \infty \text{ and} \\ &\quad \sum_{k<0} |c_k|^2 |\ell_k| < \infty\} \\ &= \{f \in H^{-1/2}(\Sigma, \mathbb{E}_\Sigma) \mid f = \sum_{k>0} c_k \varphi_k, \text{ with } \sum_{k>0} |c_k|^2 \ell_k^{-1} < \infty\} \\ &\quad + \{f \in H^{1/2}(\Sigma, \mathbb{E}_\Sigma) \mid f = \sum_{k<0} c_k \varphi_k, \text{ with } \sum_{k<0} |c_k|^2 |\ell_k| < \infty\} \\ &= \theta_+^+ + \theta_+^-. \end{aligned}$$

The maps \mathbf{i}_\pm here are given by inclusion maps.

Since $(\gamma^+(\text{Ker}(\mathcal{A}_+^*)), \theta_+^+)$ is a Fredholm pair, we can find a finite dimensional subspace F in θ_+^+ and a corresponding finite dimensional subspace G in θ_+^- such that we have decompositions

$$\theta_+^+ = F + F', \quad \theta_+^- = G + G'$$

with suitable closed subspaces F' and G' and that

$$F' + \mathbf{i}_+(G) \text{ and } \gamma^+(\text{Ker}(\mathcal{A}_+^*)) \text{ are transversal.}$$

When we put $F' + \mathbf{i}_+(G) = \lambda_-$ and $F + \overline{\mathbf{i}_+(G')} = \lambda_+$ we have a decomposition $\beta^+ = \lambda_+ + \lambda_-$ with Lagrangian subspaces λ_\pm and the Cauchy data space is expressed as a graph of a continuous map $\mathfrak{K} : \lambda_+ \rightarrow \lambda_-$. Then for such a Lagrangian subspace we have $\tau(\gamma^+(\text{Ker}(\mathcal{A}_+^*)))$ is the graph of the map $\tilde{\mathbf{i}}_- \circ \mathfrak{K} \circ \tilde{\mathbf{i}}_+$. Note here the maps $\tilde{\mathbf{i}}_\pm$ should be defined in a suitable way according to the choices of the subspace F and G (for example, $\tilde{\mathbf{i}}_+$ is defined as \mathbf{i}_+ on F' and \mathbf{i}_-^{-1} on $\mathbf{i}_-(G)$).

Now the original maps $\mathbf{i}_+ : \theta_+^- \rightarrow \theta_+^+$ and $\mathbf{i}_- : \theta_+^+ \rightarrow \theta_+^-$ are compact operators by Rellich's Theorem and so the new maps $\tilde{\mathbf{i}}_\pm$ are also compact.

Let us denote the orthogonal projection operator to a closed subspace E by \mathcal{P}_E . Then the difference

$$(3.3) \quad \mathcal{P}_{\mathbf{i}_+(F)+G'} - \mathcal{P}_{\tau(\text{Ker}(\mathcal{A}_+^*))}$$

is a compact operator and the difference

$$(3.4) \quad \mathcal{P}_{\mathbf{i}_+(F)+G'} - \mathcal{P}_{\theta_+^-}$$

is a finite rank operator.

By the definition of the map τ we have $\gamma^-(\mathfrak{D}_0) = \tau(\gamma^+(\text{Ker}(\mathcal{A}_+^*)))$, and (3.3) and (3.4) imply that the difference of the orthogonal projection operators onto the subspaces $\gamma^-(\mathfrak{D}_{APS}^0) = \theta_+^-$ and $\gamma^-(\mathfrak{D}_0)$ is a compact operator.

So this gives us a proof of Theorem (2.7).

4. CAUCHY DATA SPACES AND HÖRMANDER INDEX

Let $L_2(\Sigma) = L_+ + L_-$ be the polarization by L_\pm , where L_+ is the L_2 -completion of the space spanned by $\{\varphi_k\}_{k>0}$ and L_- is the L_2 -completion of the space spanned by $\{\varphi_k\}_{k<0}$. Then by applying above arguments to the two pairs $(\beta^+ = \beta_-^+ + \beta_+^+, L_2(\Sigma) = L_- + L_+)$ and $(\beta^- = \beta_-^- + \beta_+^-, L_2(\Sigma) = L_- + L_+)$ of polarized symplectic Hilbert spaces we have four Lagrangian subspaces

$$\gamma^\pm(\text{Ker}(\mathcal{A}_\pm^*)) \cap L_2(\Sigma), L_\pm$$

of $L_2(\Sigma)$ which satisfy following properties **(h1)** and **(h2)**:

(h1) : $\gamma^\pm(\text{Ker}(\mathcal{A}_\pm^*)) \cap L_2(\Sigma)$ and L_\mp are Fredholm pairs,

(h2) : $\gamma^\pm(\text{Ker}(\mathcal{A}_\pm^*)) \cap L_2(\Sigma) = U_\pm(L_\pm)$, where U_\pm are unitary operators of the form $\text{Id} + \text{compact operator}$.

Here we identify $L_2(\Sigma) \cong L_+ \otimes \mathbb{C}$.

Now we can define the Hörmander index

$$(4.1) \quad \sigma_H(\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma), L_+; \gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma), L_-)$$

of these four Lagrangian subspaces. Then its absolute value will express an *asymmetry* of solution spaces of the operator A under the decomposition of M along a hypersurface Σ . So, if there is a *symmetry* among these four Lagrangian subspaces, the value must vanish. In fact

Proposition 4.1. *Assume that $\sigma(\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)) = \gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma)$, then the Hörmander index of these four Lagrangian subspaces vanishes:*

$$\sigma_H(\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma), L_+; \gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma), L_-) = 0.$$

Proof. First we assume that $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$ and L_- are transversal. Then the space $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$ is written as a graph of a compact operator $T : L_+ \rightarrow L_-$ such that $\sigma \circ T$ is a selfadjoint operator on L_+ and the space $\gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma)$ is also written as a graph of the map $-\sigma \circ T \circ \sigma$. These imply that the curve of Lagrangian subspaces given by the graphs of $\{-t \cdot \sigma \circ T \circ \sigma\}_{0 \leq t \leq 1}$ is always transversal to both of $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$ and L_+ . This curve is connecting $\gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma)$ and L_- . Hence we have

$$\sigma_H(\gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma), L_-; \gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma), L_+) = 0.$$

If $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$ and L_- are not transversal, then we decompose the Lagrangian subspace $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$ into the orthogonal sum

$$\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma) = \ell_0 + \nu,$$

where $\ell_0 = (\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)) \cap L_-$ and ν is the orthogonal complement of ℓ_0 in $\gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma)$. Also we decompose $L_- = \ell_0 + (L_- \cap \ell_0^\perp) = \ell_0 + \ell_-$ and $L_+ = \sigma(\ell_0) + (L_+ \cap \sigma(\ell_0)^\perp) = \sigma(\ell_0) + \ell_+$. Now we have

$$\begin{aligned} & \sigma_H(\gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma), L_-; \gamma^+(\text{Ker}(\mathcal{A}_+^*)) \cap L_2(\Sigma), L_+) \\ &= \sigma_H(\sigma(\ell_0), \ell_0; \ell_0, \sigma(\ell_0)) + \sigma_H(\sigma(\nu), \ell_-; \nu, \ell_+) = 0, \end{aligned}$$

by applying the first arguments to the second term.

Note that $\gamma^-(\text{Ker}(\mathcal{A}_-^*)) \cap L_2(\Sigma) = \sigma(\ell_0) + \sigma(\nu)$ is an orthogonal decomposition and the vanishing of the first term follows from a skew-symmetric property of the Hörmander index. \square

5. A SPLITTING FORMULA OF A SPECTRAL FLOW

First we state a splitting formula for a spectral flow when we decompose a manifold into two components. Then we give another form of it by replacing the boundary condition with Atiyah-Patodi-Singer condition.

Let $\{C_t\}_{t \in [0,1]}$ be a continuous family of symmetric bundle maps of \mathbb{E} and we assume that each of the operator in the family $\{A + C_t\}$ satisfies the conditions **(a1')** and **(a2)** where **(a1')** is:

(a1'): There exists an $\epsilon_0 > 0$ such that for any $|s| < \epsilon_0$ and any $t \in [0,1]$ the operators $A + C_t + s$ satisfy the *unique continuation property* with respect to the hypersurface Σ :

$$(5.1) \quad \text{Ker}(\mathcal{A}_\pm^* + C_t + s) \cap \mathfrak{D}_m^\pm = \{0\}.$$

Here C_t is regarded as a bounded selfadjoint operator on $L_2(M, \mathbb{E})$.

Now we have continuous families of Cauchy data spaces $\gamma^\pm(\text{Ker}(\mathcal{A}_\pm^* + C_t))$ (C_t should be considered as acting on the space \mathfrak{D}_M^\pm respectively, and both of which are invariant under this action). The splitting formula is stated as follows:

Theorem 5.1. ([FO])

$$(5.2) \quad \mathbf{Sf}(\{A + C_t\}) = \mathbf{Sf}(\{A_{\mathfrak{D}_0} + C_t\}) + \mathbf{Sf}(\{A_{\mathfrak{D}_1} + C_t\}),$$

where

$$\mathfrak{D}_0 = \{f \in H^1(M_-, \mathbb{E}|_{M_-}) \mid \exists \tilde{f} \in H^1(M, \mathbb{E}) \text{ such that } \tilde{f}|_{M_-} = f \text{ and } (A+C_0)(\tilde{f}) = 0 \text{ on } M_+\}$$

and

$$\mathfrak{D}_1 = \{g \in H^1(M_+, \mathbb{E}|_{M_+}) \mid \exists \tilde{g} \in H^1(M, \mathbb{E}) \text{ such that } \tilde{g}|_{M_+} = g \text{ and } (A+C_1)(\tilde{g}) = 0 \text{ on } M_-\}.$$

Remark 5.2. Our proof of the general spectral flow formula bases on the property **(a1')** and **(a2)**, and by making use of the general spectral flow formula and L_2 reduction theorem we prove the splitting formula above ([FO]).

Let \mathfrak{D}_{APS}^0 be the space defined in (2.6) for A replaced by $A + C_0$ and denote by \mathfrak{D}_{APS}^1 the similar space

$$(5.3) \quad \mathfrak{D}_{APS}^1 = \left\{ f \in H^1(M_+, \mathbb{E}|_{M_+}) \mid \text{if } f|_{\Sigma} = \sum_{k \in \mathbb{Z} \setminus \{0\}} a_k \psi_k, \text{ then } a_k = 0 \text{ for } k = 1, 2, 3, \dots \right\}.$$

Note that the sections $\{\psi_k\}$ are now orthonormal eigensections of the tangential operator B_1 in the product form

$$(5.4) \quad A + C_1 = \sigma \left(\frac{\partial}{\partial u} + B_1 \right)$$

corresponding to the parameter $t = 1$ and should be chosen in such a way as noted in Remark (2.4).

We have continuous curves $\{\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_t))\}$ of Cauchy data spaces in the Fredholm-Lagrangian-Grassmannian $\mathcal{F}\Lambda_{\mathfrak{D}_0}(\beta^-) = \mathcal{F}\Lambda_{\mathfrak{D}_{APS}^0}(\beta^-)$ and $\{\gamma^+(\text{Ker}(\mathcal{A}_+^* + C_t))\}$ in $\mathcal{F}\Lambda_{\mathfrak{D}_1}(\beta^+) = \mathcal{F}\Lambda_{\mathfrak{D}_{APS}^1}(\beta^+)$.

The Hörmander index is defined for four Lagrangian subspaces $\mu_0, \mu_1, \gamma^-(\mathfrak{D}_0)$ and $\gamma^-(\mathfrak{D}_{APS}^0)$, where $\mu_i \in \mathcal{F}\Lambda_{\mathfrak{D}_0}(\beta^-) = \mathcal{F}\Lambda_{\mathfrak{D}_{APS}^0}(\beta^-)$, also defined for $\nu_0, \nu_1, \gamma^+(\mathfrak{D}_1)$ and $\gamma^+(\mathfrak{D}_{APS}^1)$ ($\nu_i \in \mathcal{F}\Lambda_{\mathfrak{D}_1}(\beta^+) = \mathcal{F}\Lambda_{\mathfrak{D}_{APS}^1}(\beta^+)$), as noted in Remark (2.8).

Since

$$\begin{aligned} & \mathbf{Sf}(\{A_{\mathfrak{D}_0} + C_t\}) \\ &= \mathbf{Mas}(\{\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_t))\}, \gamma^-(\mathfrak{D}_0)) \\ &= \mathbf{Mas}(\{\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_t))\}, \gamma^-(\mathfrak{D}_{APS}^0)) \\ &\quad + \sigma_H(\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_0)), \gamma^-(\text{Ker}(\mathcal{A}_-^* + C_1)); \gamma^-(\mathfrak{D}_0), \gamma^-(\mathfrak{D}_{APS}^0)) \end{aligned}$$

we have

Theorem 5.3.

$$\begin{aligned} & \mathbf{Sf}(\{A + C_t\}) \\ &= \mathbf{Mas}(\{\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_t))\}, \gamma^-(\mathfrak{D}_{APS}^0)) \\ &\quad + \sigma_H(\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_0)), \gamma^-(\text{Ker}(\mathcal{A}_-^* + C_1)); \gamma^-(\mathfrak{D}_0), \gamma^-(\mathfrak{D}_{APS}^0)) \\ &+ \mathbf{Mas}(\{\gamma^+(\text{Ker}(\mathcal{A}_+^* + C_t))\}, \gamma^+(\mathfrak{D}_{APS}^1)) \\ &\quad + \sigma_H(\gamma^+(\text{Ker}(\mathcal{A}_+^* + C_0)), \gamma^+(\text{Ker}(\mathcal{A}_+^* + C_1)); \gamma^+(\mathfrak{D}_1), \gamma^+(\mathfrak{D}_{APS}^1)) \\ &= \mathbf{Sf}(\{A_{\mathfrak{D}_{APS}^0} + C_t\}) + \mathbf{Sf}(\{A_{\mathfrak{D}_{APS}^1} + C_t\}) \\ &\quad + \sigma_H(\gamma^-(\text{Ker}(\mathcal{A}_-^* + C_0)), \gamma^-(\text{Ker}(\mathcal{A}_-^* + C_1)); \gamma^-(\mathfrak{D}_0), \gamma^-(\mathfrak{D}_{APS}^0)) \\ &\quad + \sigma_H(\gamma^+(\text{Ker}(\mathcal{A}_+^* + C_0)), \gamma^+(\text{Ker}(\mathcal{A}_+^* + C_1)); \gamma^+(\mathfrak{D}_1), \gamma^+(\mathfrak{D}_{APS}^1)). \end{aligned}$$

Corollary 5.4. *If the family $\{A + C_t\}$ is a loop, i.e., $C_0 = C_1$, then we have*

$$(5.5) \quad \mathbf{Sf}(\{A + C_t\}) = \mathbf{Sf}(\{A_{\mathfrak{D}_{APS}^0} + C_t\}) + \mathbf{Sf}(\{A_{\mathfrak{D}_{APS}^1} + C_t\}).$$

Remark 5.5. Although it holds the spectral flow formula expressed in terms of the Maslov index of Cauchy data spaces under the assumption **(a1')**, it would not be clear whether the splitting formulas of the spectral flow like above formulas hold always without the second assumption **(a2)**. Such assumptions are fit to the framework of the symplectic

Hilbert space theory, after once the spaces β^\pm are determined. However it would be expected that generalizations of splitting formula of spectral flow and the index similar to (4.1) without the assumption **(a2)** would be carried out through a further analysis of the pseudo-differential operator theory including the Calderón projector and the operator $\mathcal{P}_{\tau(\text{Ker}(\mathcal{A}_+^*))}$.

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